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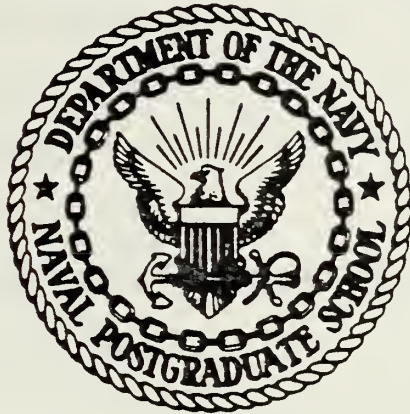






# NAVAL POSTGRADUATE SCHOOL

## Monterey, California



# THESIS

TIME INTEGRATION OF UNSTEADY TRANSONIC FLOW  
TO A STEADY STATE SOLUTION BY  
THE FINITE ELEMENT METHOD

by

Raymond John Nichols Jr.

March 1977

Thesis Advisor:

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Time Integration of Unsteady Transonic Flow  
to a Steady State Solution by  
the Finite Element Method

by

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Lieutenant, United States Navy  
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Submitted in partial fulfillment of the  
requirements for the degree of

MASTER OF SCIENCE IN AERONAUTICAL ENGINEERING

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## ABSTRACT

A finite element method was applied to the unsteady transonic small disturbance equation and integrated until the solution converged to the steady state for a thin non-lifting airfoil. The method of weighted residuals was used to formulate the finite element equations, and Houbolt's method of central differencing in time was used to integrate these equations.

A secondary investigation applied the steady transonic small disturbance equations to a converging-diverging nozzle.



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## I. INTRODUCTION

Transonic inviscid flows past a smooth airfoil may be expressed in terms of the velocity potential  $\phi$  satisfying the transonic small disturbance equation,

$$(1 - M_\infty^2 - M_\infty^2(\gamma + 1)\phi_x)\phi_{xx} + \phi_{yy} = 0 \quad (1)$$

This equation presents two major difficulties, 1) it is non-linear and 2) it is of the mixed hyperbolic-elliptic type. Analytical solutions to non-linear equations are difficult to obtain. One must normally resort to numerical methods. When the coefficient  $(1 - M_\infty^2 - M_\infty^2(\gamma + 1)\phi_x)$  in equation 1 is negative, the flow is supersonic and the equation is called hyperbolic; otherwise the flow is subsonic and the equation is elliptic. The forms of the two solutions are fundamentally different. Hyperbolic equations admit both discontinuities, which propagate only in characteristic directions, and the presence of shock fronts. Elliptic equations, on the other hand, require that the dependent variables and their derivatives be continuous and that a change in any part of the flow field affects every other part. Many non-linear elliptic equations are solved by appropriate relaxation iteration techniques by casting the equation in Poisson's form with the non-linearity acting as the driving force. Solutions to hyperbolic equations are usually obtained by the method of characteristics or by



finite difference marching techniques which use an artificial viscosity to represent the average jump conditions across the shock wave. In mixed supersonic and subsonic flows, normal numerical procedures break down because the boundary between the two regions is not known a priori.

Finite element numerical techniques have evolved as a powerful tool in obtaining approximate solutions to a wide variety of engineering problems, particularly ones with Neumann-Dirichlet boundary conditions, i.e., elliptical problems. They offer several outstanding advantages. Some of these are:

- 1) Non-homogeneous problems may be treated with relative simplicity.
- 2) Complex geometries may be modeled with relative ease since the elements can be graded in size and shape to follow boundaries of arbitrary shape.
- 3) Once the finite element model has been established, a variety of problems can be solved by supplying the computer with the appropriate data.

Chan and Brashears [Ref. 5] developed a finite element computer program for steady transonic flow over a non-lifting airfoil. This program uses the least squares method of weighted residuals to approximate equation 1 by a system of algebraic equations, and assembles the equations in a special way to account for the hyperbolic region of flow. This technique prevents the influence of downwind nodes from propagating upstream in the supersonic region.

The purpose of this thesis is to investigate the possibility of speeding the convergence of a solution by transforming the steady transonic equation to the unsteady equation



and integrating over time until the time dependent terms vanish and to extend the program of Ref. [5] to the transonic region of a converging-diverging nozzle.





## II. DISCUSSION OF THE FINITE ELEMENT APPROACH

In a continuum problem of any dimension, the field variable, whether it is velocity potential, velocity, temperature, displacement or some other quantity, possesses infinitely many values because it is a function of each generic point in the solution region. Consequently, the problem is one of an infinite number of unknowns. The finite element approach subdivides the solution domain into a finite number of subdomains called elements and expresses the unknown field variable in terms of assumed approximating functions within each element. The approximating functions are sometimes called interpolating functions and are defined in terms of the values of the field variables at specified points called nodes or nodal points. Nodes usually lie on the element boundaries where adjacent elements are considered to be connected. In addition to boundary nodes, an element may also contain interior nodes. The nodal values of the field variable and the interpolating function for the elements completely define the behavior of the field variable within the elements. Once these unknowns are found, the interpolating functions define the field variable throughout the assemblage of the element.

Clearly, the nature of the solution and the degree of accuracy of the approximation depends not only on the number and size of the elements used but also on the interpolating



functions selected. Interpolating functions may not be chosen arbitrarily because certain compatibility conditions must be satisfied. Often such functions are selected so that the field variable or its derivatives are continuous across adjoining element boundaries. Once the problem is formulated in terms of individual elements, the contributions of each element may be assembled to define the entire solution domain. This means, for example, that if we are treating a problem in stress analysis we can find the force-displacement or stiffness characteristics of each element and then assemble the elements to determine the stiffness of the whole structure. Finite element solutions are not, of course, restricted to structures problems, but the matrix of equations defined by the interpolating functions and the nodal field variables is still referred to as the stiffness matrix regardless of the field variable in the problem.

Solutions to continuous problems by the finite element approach always follow a systematic step-by-step process. This process is completely general to the finite element method and it is outlined below. [Ref. 4]

#### 1. Discretize the continuum.

The first step is to divide the solution domain into elements. A variety of element shapes may be used and one or more different element shapes may be used in the same region. The type and number of the elements used in a given problem are a matter of engineering judgement.



## 2. Select the interpolating functions.

The next step is to choose the type of interpolating function to represent the variation of the field variable over the element. The field variable may be a scalar, a vector, or a higher order tensor. Often polynomials are selected as interpolating functions for they are easy to integrate and differentiate. The degree of polynomial chosen depends on the number of nodes and the nature and number of the unknowns and the continuity requirements imposed at the nodes and the element boundaries. The unknown quantities at the nodes may be assigned to the field variable and their derivatives.

## 3. Find the element properties.

After the elements and their interpolating functions have been selected, the matrix of algebraic equations which express the properties of the individual elements must be determined. Several methods are available for this task. These are: the direct approach, the variational approach, the method of weighted residuals and the energy balance approach. Reference [4] is a good source of information on the various techniques.

## 4. Assemble the element properties to obtain the system equations.

To find the properties of the over-all system, the matrix equations expressing the behavior of each element must be



added to the matrix equation of all other elements. The basis for this assembly procedure stems from the fact that connecting elements have common nodes and the field variable must be the same for each element sharing that node.

At this point the boundary conditions for the system of equations must be accounted for and the system of equations must be modified before it is ready for solution.

5. Solve the system of equations.

The assembly process of step 4 produces a set of simultaneous equations which can be solved to obtain the unknown field variables. Linear equations have a number of standard solution techniques readily available, but solutions to non-linear equations are more difficult to obtain.

6. Make additional computations if desired.

Important parameters, such as pressure coefficient in aerodynamic problems, may now be calculated from the values of the field variables.





### III. THEORY AND BASIC EQUATIONS

#### A. STEADY TRANSONIC FLOW

Chan et al. [Ref. 5] developed an algorithm to analyze steady transonic flow over non-lifting thin airfoils. Boundary layer effects were ignored and the imbedded shock wave was assumed to be weak. These assumptions are consistent with transonic small disturbance theory which can be expressed mathematically by the following expressions.

$$(1 - M_{\infty}^2 - M_{\infty}^2(\gamma + 1)\phi_x)\phi_{xx} + \phi_{yy} = 0 \quad (1)$$

Boundary conditions -

$$\nabla \cdot \phi = 0 \quad \text{at infinity} \quad (2)$$

$$v = (1 + u)dg/dx \quad \text{on the airfoil} \quad (3)$$

$$u = 0 \quad \text{on the line of symmetry} \quad (4)$$

where  $g(x,y)$  defines the airfoil and  $dg/dx$  describes the airfoil slope.

The above expressions appear in their dimensionless form where  $\phi$  = perturbed velocity potential and the perturbed velocity components in the  $x$  and  $y$  directions are respectively defined as

$$u = \phi_x$$

$$v = \phi_y$$



$M_\infty$  = freestream Mach number and  $\gamma$  = ratio of specific heats which for air is taken to be 1.4. The physical coordinates  $x'$  and  $y'$  and the velocity potential  $\phi'$  are related to the dimensionless quantities by

$$x = x'/c, \quad y = y'/c, \quad \phi = \phi'/cU_\infty$$

where  $c$  is the chordlength of the airfoil and  $U_\infty$  is the free-stream velocity.

Once the flowfield solution is determined in terms of the perturbed velocities, the secondary unknowns are subsequently calculated. These include:

$$a = \left[ \frac{\gamma-1}{2} (U_\infty^2 - V^2) + \left( \frac{U_\infty}{M} \right)^2 \right]^{\frac{1}{2}} \quad (5)$$

$$M = \frac{V}{a} \quad (6)$$

$$\frac{P}{P_0} = \frac{1}{\left[ 1 + \frac{\gamma-1}{2} M^2 \right]^{\gamma/(\gamma+1)}} \quad (7)$$

$$C_p = - \left[ \frac{2u}{U_\infty} + (1-M_\infty^2) \frac{u^2}{U_\infty^2} + \frac{v^2}{U_\infty^2} \right] \approx - \frac{2u}{U_\infty} \quad (8)$$

In the above,  $U_\infty = 1$  is the normalized freestream velocity,  $a$  = local sound speed,  $p$  = local static pressure,  $M$  = local Mach number,  $V$  = total velocity,  $P_0$  = stagnation pressure and  $C_p$  = pressure coefficient.

## B. UNSTEADY TRANSONIC FLOW

Unsteady transonic inviscid flow may be expressed in terms of the velocity potential  $\phi(x,y,t)$  to a first order approximation by



$$(1 - M_\infty^2 - M_\infty^2(\gamma + 1)\phi_x)\phi_{xx} + \phi_{yy} - 2M_\infty^2\phi_t - M_\infty^2\phi_{tt} = 0 \quad (9)$$

which has the same non-linear coefficient retained for steady transonic flow in Equation 1.

Boundary conditions require that the disturbances vanish far from the airfoil,

$$\phi_x = 0 \quad \phi_y = 0 \quad \text{at infinity} \quad (10)$$

and that the flow remain attached to the body. Let  $B(x,y,t) = 0$  define the body at any instant. The surface tangency restraint may now be expressed by the substantial derivative  $DB/DT$  vanishing.

$$DB/DT = B_t + (1 + \phi_x)B_x + \phi_y B_y \quad (11)$$

If the body remains stationary  $B_t = 0$ , and the tangency condition becomes the same as in steady flow, namely

$$v = (1 + u)dg/dx \quad (12)$$

where  $dg/dx$  represents the airfoil slope.

The pressure coefficient for isentropic unsteady compressible flow is defined by

$$C_p = \frac{2}{\gamma M_\infty^2} \left\{ \left[ 1 - \frac{\gamma-1}{2} M_\infty^2 (2\phi_x + 2\phi_t + \phi_x^2 + \phi_y^2) \right]^{\gamma/(\gamma-1)} - 1 \right\} \quad (13)$$

Expanding by the binomial expansion and retaining only the first order terms gives,

$$C_p = -2\phi_x - 2\phi_t \quad (14)$$



#### IV. FINITE ELEMENT FORMULATION

The method of weighted residuals is a technique for approximating solutions to linear or non-linear partial differential equations and it is the basis for the finite element formulation of the transonic small disturbance equation (Equation 1). This procedure involves assuming the general functional behavior of the field variable which would approximately satisfy the basic equation and boundary conditions. Substituting this approximation into the original differential equation results in some error called a residual. A system of algebraic equations results when a weighted average of the residual is forced to vanish as it is averaged over the entire domain.

##### A. STEADY FLOW

The approximate solution to equation 1 is assumed to be

$$\hat{\phi} = N_i \phi_i \quad (15)$$

where  $N_i$  are the interpolating functions which exhibit the behavior of equation 1 and  $\phi_i$  are the undetermined parameters at the nodal points.

When  $\hat{\phi}$  is substituted into equation 1, the resulting residual is

$$R = [1 - M_\infty^2 - M_\infty^2(\gamma + 1)N_{k,x}\phi_k]N_{j,xx} + N_{j,yy} \quad (16)$$





The weighted average is determined by multiplying the residual  $R$  by  $m$  linearly independent weighting functions  $W_i$  and integrating over the elemental domain. Forcing this residual to vanish yields,

$$\iint W_i R dA = 0 \quad i = 1 \text{ to } m$$

Chan et al. [Ref. 5] found that when the weighting function  $W_i$  for equation 1 is chosen to be  $\partial R / \partial \phi_i$  the resulting matrix is positive definite and well conditioned. This choice of weighting functions is referred to as the method of least squares because it is equivalent to minimizing the square of the residuals summed over the domain with respect to the undetermined parameters. That is,

$$\begin{aligned} \chi &= \iint R^2 dA \\ \partial \chi / \partial \phi_i &= \iint \partial R / \partial \phi_i R dA = 0 \end{aligned} \quad (17)$$

Integrating over the domain produces the system of algebraic equations

$$S_{ij} \phi_j = 0 \quad (18)$$

where the elemental matrix  $S_{ij}$  is defined as

$$S_{ij} = \iint Q_j P_i dA \quad (19)$$

With  $Q_j$  and  $P_i$  equal to

$$Q_j = [1 - M_\infty^2 - M_\infty^2(\gamma + 1)N_{k,x}\phi_k] N_{j,xx} + N_{j,yy}$$

$$P_i = Q_i - M_\infty^2(1 + \gamma)N_{k,xx}\phi_k N_{i,x}$$



## B. UNSTEADY FLOW

Development of the unsteady flow finite element equations is similar to the procedure used to formulate the finite element equations for steady transonic flow. The least squares method of weighted residuals is again used but  $\phi$  is now a function of time as well as the spatial coordinates  $x$  and  $y$ .

The approximate solution has the form,

$$\hat{\phi} = N_i(x,y)\phi_i(t) \quad (20)$$

Substituting  $\hat{\phi}$  in the unsteady transonic small disturbance equation, the weighted residual becomes,

$$\chi = \iint (R_1 + R_2 + R_3)^2 dA \quad (21)$$

where

$$R_1 = \{ [1 - M_\infty^2 - M_\infty^2(\gamma + 1)\phi_k N_{k,x}] N_{j,xx} + N_{j,yy} \} \phi_j$$

$$R_2 = -2M_\infty^2 N_{j,x} \dot{\phi}_j$$

$$R_3 = -M_\infty^2 N_j \ddot{\phi}_j$$

Expanding equation 21 yields

$$\chi = \iint [R_1^2 + R_2^2 + R_3^2 + 2R_1R_2 + 2R_2R_3] dA \quad (22)$$

and minimizing with respect to the undetermined parameters  $\phi_i$  the following system of algebraic equations result,

$$\phi_j = 0 = 2 \iint \partial R_i / \partial \phi_j [R_1 + R_2 + R_3] dA$$

$$\partial R_i / \partial \phi_j = P_i$$



where  $P_i$  has been previously defined in the steady finite element formulation

The above equation may be rewritten in the form

$$S_{ij}\phi_j + SC_{ij}\dot{\phi}_j + SM_{ij}\ddot{\phi}_j = 0 \quad (23)$$

The stiffness matrix  $S_{ij}$  is the same as that developed for the steady transonic equation and the damping ( $SC_{ij}$ ) and mass ( $SM_{ij}$ ) matrices are defined below.

$$SC_{ij} = -\iint M_\infty^2 N_{j,x} P_i dA \quad (24)$$

$$SM_{ij} = -\iint M_\infty^2 N_j P_i dA \quad (25)$$



## V. ELEMENT DESCRIPTION AND ASSEMBLY OF EQUATIONS

### A. ELEMENT DESCRIPTION

The basic element used in the finite element program is the non-conforming cubic triangular element developed by Bazeley et al. [Ref. 2]. Also used in the program are the quadrilateral and trapezoidal elements constructed from these triangular elements. These three types of elements can be mixed and used freely in the entire flow region except that only trapezoids should be used to cover the supersonic and mixed region in order to enact the special assembly procedures required by the hyperbolic equation which describes the flow in that region.

The basic triangular element is shown in Fig. 1, which at each vertex has the velocity potential and the velocity components as the undetermined parameters. This type of element was chosen because both Dirichlet and Neumann boundary conditions can be imposed with equal convenience. In addition, because velocity components are used as primary unknowns secondary parameters, such as Mach number and pressure coefficient can be calculated directly without resorting to numerical differentiation, which would produce additional errors.

In the element, the approximate solution is assumed as

$$\hat{\phi} = N_i \phi_i \quad (i = 1 \text{ to } 9)$$





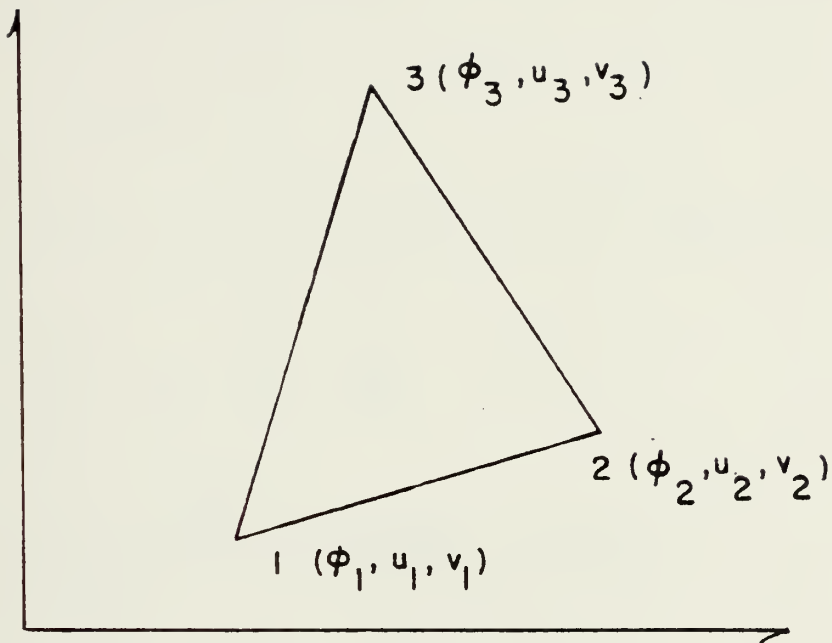


Figure 1 - Triangular Element



In which  $\phi_i$ 's are the nine undetermined parameters of  $\phi$  and  $N_i$  are the interpolation functions which are expressed in terms of the area coordinates.

The shape or interpolating functions and their first and second derivatives are defined below.

Letting,

$$a_i = x_j y_k - x_k y_j$$

$$b_i = y_j - y_k$$

$$c_i = x_k - x_j$$

$$\Delta = \text{area of triangle } 1-2-3 = (b_j c_k - b_k c_j)/2$$

$$\alpha = 0.5 (c_k - c_j)$$

$$\beta = 0.5 (b_j - b_k)$$

$$H = \zeta_k \zeta_j \zeta_k$$

$$H_x = b_i \zeta_j \zeta_k + b_j \zeta_k \zeta_i + b_k \zeta_i \zeta_j$$

$$H_y = c_i \zeta_j \zeta_k + c_j \zeta_k \zeta_i + c_k \zeta_i \zeta_j$$

$$H_{xx} = 2(\zeta_k b_j b_k + \zeta_j b_k b_k + \zeta_k b_i b_j)$$

$$H_{yy} = 2(\zeta_i c_j c_k + \zeta_j c_k c_i + \zeta_k c_i c_j)$$

with  $i = (1,2,3), k = (3,1,2)$  then one has for  $l = (1,4,7),$   
 $i = (1,2,3)$

$$N_1 = \zeta_i^2 (3 - 2\zeta_i) + 2H$$

$$N_{1,x} = [6b_i \zeta_i (1 - \zeta_i) + 2H_x]/2\Delta$$

$$N_{1,y} = [6c_i \zeta_i (1 - \zeta_i) + 2H_y]/2\Delta$$



$$N_{1,xx} = [6b_i^2 (1 - 2\zeta_i) + 2H_{xx}] / (2\Delta)^2$$

$$N_{1,yy} = [6c_i^2 (1 - 2\zeta_i) + 2H_{yy}] / (2\Delta)^2$$

for  $1 = (2, 5, 8)$ ,  $1 = (1, 2, 3)$

$$N_1 = \zeta_i^2 (c_k \zeta_j - c_j \zeta_k) + \alpha H$$

$$N_{1,x} = [2b_i \zeta_i (c_k \zeta_j - c_j \zeta_k) + 2\zeta_i^2 + \alpha H_x] / 2\Delta$$

$$N_{1,y} = [2c_i \zeta_i (c_k \zeta_j - c_j \zeta_k) + \alpha H_y] / 2\Delta$$

$$N_{1,xx} = [2b_i^2 (c_k \zeta_j - c_j \zeta_k) + 4b_i (2\Delta) \zeta_i + \alpha H_{xx}] / (2\Delta)^2$$

$$N_{1,yy} = [2c_i^2 (c_k \zeta_j - c_j \zeta_k) + \alpha H_{yy}] / (2\Delta)^2$$

for  $1 = (3, 6, 9)$ ,  $1 = (1, 2, 3)$

$$N_1 = \zeta_i^2 (b_j \zeta_k - b_k \zeta_j) + \beta H$$

$$N_{1,x} = [2b_i \zeta_i (b_j \zeta_k - b_k \zeta_j) + \beta H_x] / 2\Delta$$

$$N_{1,y} = [2c_i \zeta_i (b_j \zeta_k - b_k \zeta_j) + 2\zeta_i^2 + \beta H_y] / 2\Delta$$

$$N_{1,xx} = [2b_i^2 (b_j \zeta_k - b_k \zeta_j) + \beta H_{xx}] / (2\Delta)^2$$

$$N_{1,yy} = [2c_i^2 (b_j \zeta_k - b_k \zeta_j) + 4c_i (2\Delta) \zeta_i + \beta H_{yy}] / (2\Delta)^2.$$



Quadrilateral and trapezoidal elements as shown in Fig. 2 are also used in the present program, the former is used in the subsonic region and the latter in the mixed and supersonic region. The element matrix for the quadrilateral element is obtained by combining appropriately the matrices for two triangles, while the matrix for trapezoidal element is obtained by averaging contributions from two left-running and two right-running triangles. The averaging process is designed to remove the bias effects inherent in the quadrilaterals used.

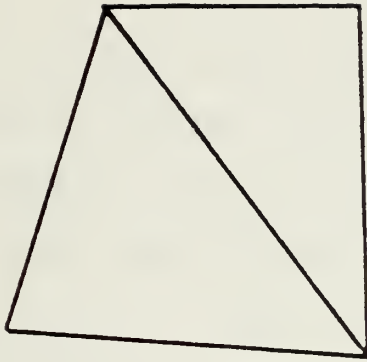
## B. ASSEMBLY OF EQUATIONS

Straightforward application of the finite element assembly technique to transonic flows would fail (the solution diverges) because this would allow disturbances to propagate upwind in the supersonic region of flow where the governing equation is hyperbolic. Hyperbolic equations have a time-like dependency in that the solution at the downwind station is affected by the upwind station but not vice-versa. Assembly techniques for a transonic flow finite element program must take into account this time-like dependency. If the x-axis is taken as a time-like direction in the supersonic region, the element matrix may be assembled in a way similar to a backward finite difference operator, which has been successful in solving hyperbolic equations.

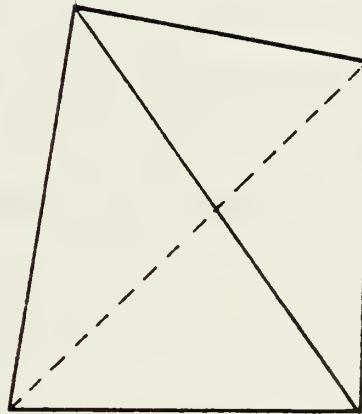
Consider the rectangular element sketched below with the upwind station I and the downwind station II, each having two nodal points with the element type chosen. The element matrices can be constructed in the usual manner.







a. Quadrilateral  
Element



b. Trapezoidal  
Element

Figure 2 - Quadrilateral and Trapezoidal Elements



However, before assembling the element matrix into the system matrix the non-linear coefficient of equation 1 is evaluated.

$$C = 1 - M_{\infty}^2 - M_{\infty}^2 (\gamma + 1) u$$

The sign of the coefficient being positive, zero, or negative defines the equation as elliptic, parabolic, or hyperbolic. If C is non-positive for all nodes in the element, the rows representing the improper downwind influence on the solution at an upwind station are ignored during assembly. This feature is automatically applied in the program requiring only a little care in arrangement of the nodes of the element. In the anticipated supersonic region, element node points should be arranged in the order as indicated in figure 3, starting with the upper left corner and proceeding in the counter-clockwise direction. In the elliptic region, i.e., where the coefficient is positive, no special assembly technique is invoked.

### C. ITERATIVE PROCEDURES

With the equations assembled and the proper boundary conditions imposed, the system of non-linear algebraic equations is solved by iterative procedures in the form

$$S_{ij}(\tilde{\phi})\phi_j^{(n)} = 1_i \quad (23)$$

to solve for the solution in the  $n^{\text{th}}$  iteration. The function  $\tilde{\phi}$  is defined as

$$\tilde{\phi} = \theta\phi^{(n-1)} + (1-\theta)\phi^{(n-1)}$$



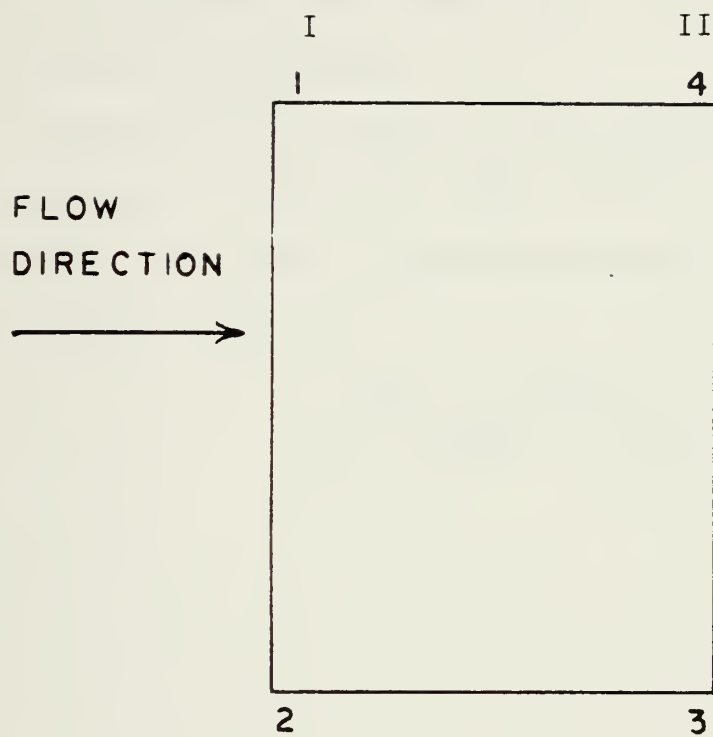


Figure 3 - Nodal Arrangement for Supersonic Region



in which the under-relaxation factor  $\theta$  is in the range  $0 < \theta \leq 1$ . For subsonic flow  $\theta = 1$ , which is simply a successive approximation, yields good results, but it is necessary to under-relax somewhat with  $\theta$  approximately .5 for supercritical flow. Generally, a smaller relaxation factor will make the solution more stable but it will tend to slow down the rate of convergence.

Equation 23 is subject to the convergence criterion that the change in local Mach number between two successive iterations is less than a prescribed value  $\epsilon$  at all nodes in the flow field. That is,

$$\left| \frac{M^{(n)} - M^{(n-1)}}{M^{(n)}} \right| \leq \epsilon .$$





## VI. INTEGRATION OF UNSTEADY FINITE ELEMENT EQUATIONS

The unsteady transonic small disturbance equation (equation 9) when suitably reduced to a finite element approximation appears in the form,

$$S_{ij}\phi_j + SC_{ij}\dot{\phi}_j + SM_{ij}\ddot{\phi}_j = R_j \quad (26)$$

This equation is analogous to a damped spring mass system, hence  $S_{ij}$ ,  $SC_{ij}$ , and  $SM_{ij}$  are respectively referred to as the stiffness, damping, and mass matrices.

Mathematically, equation 26 represents a system of second order differential equations with constant coefficients, which can be solved by standard numerical procedures for differential equations, such as Runge-Kutta or Milne methods. However, this would be a very costly technique if the coefficient matrices are very large. In practical finite element analysis there are a few effective methods which take advantage of the banded matrices usually encountered in finite elements. One such method is the direct numerical integration method.

Direct integration involves a numerical step-by-step procedure aimed at satisfying equation 26 only at discrete time intervals  $\Delta t$  apart and not over all time  $t$ . Conceptually, direct integration is a finite element method in space and a finite difference method in time. Examples of direct integration are the central difference method, Houbolt integration and the Wilson method. The first two schemes are finite



difference schemes whereas the latter is a linear acceleration method. Linear acceleration integration assumes a linear variation of acceleration from time  $t$  to time  $t + \Delta t$ .

Central differencing can be very effective in the solution of many dynamic problems especially those that involve a large system of equations. However, this method is unstable for all time steps larger than a critical time step.

Houbolt integration is an implicit finite differencing method related to central differencing, only it has the advantage of being stable for all TIME STEPS.

The Houbolt method was used to integrate the unsteady finite element equations because of this stability.

Houbolt integration uses the following finite difference expansions:

$$\phi_{i,t+\Delta t} = [2\phi_{i,t+\Delta t} - 5\phi_{i,t} + 4\phi_{i,t-\Delta t} - \phi_{i,t-2\Delta t}]/\Delta t^2$$

$$\phi_{i,t+\Delta t} = [11\phi_{i,t+\Delta t} - 18\phi_{i,t} + 9\phi_{i,t-\Delta t} - 2\phi_{i,t-2\Delta t}]/6\Delta t$$

which are two backward-difference formulas with errors of order  $(\Delta t)^2$ .

The solution of  $\phi_{i,t+\Delta t}$  must satisfy equation 26 and at time  $t+\Delta t$  equation 26 becomes

$$S_{ij}\phi_{j,t+\Delta t} + SC_{ij}\dot{\phi}_{j,t+\Delta t} + SM_{ij}\ddot{\phi}_{j,t+\Delta t} = 0$$

Substituting the finite difference formulas for  $\phi_{j,t+\Delta t}$  and rearranging all known vectors on the right hand side, the solution for  $\phi_{j,t+\Delta t}$  is obtained, namely:



$$(a_0^{SM_{ij}} + a_1^{SC_{ij}} + S_{ij})\phi_{j,t+\Delta t} = R_{j,t+\Delta t} \quad (27)$$

$$+ (a_2^{SM_{ij}} + a_3^{SC_{ij}})\phi_{j,t} + (a_4^{SM_{ij}} + a_5^{SC_{ij}})\phi_{j,t-2\Delta t}$$

Where the constant integration coefficients are:

$$a_0 = 2/\Delta t^2$$

$$a_1 = 11/6\Delta t$$

$$a_2 = 5/\Delta t$$

$$a_3 = 3/\Delta t$$

$$a_4 = 2a_0$$

$$a_5 = -a_3/2$$

$$a_6 = a_0/2$$

$$a_7 = a_3/9$$

Equation 27 may be written as

$$SE_{ij}\phi_{j,t+\Delta t} = RE_j \quad (28)$$

where the effective stiffness matrix  $SE_{ij}$  and the effective load vector  $RE_j$  are defined as:

$$SE_{ij} = S_{ij} + a_0^{SM_{ij}} + a_1^{SC_{ij}}$$

$$RE_j = SM_{ij}(a_2\phi_{j,t} + a_4\phi_{j,t-\Delta t} + a_6\phi_{j,t-2\Delta t}) \\ + SC_{ij}(a_3\phi_{j,t} + a_5\phi_{j,t-\Delta t} + a_7\phi_{j,t-2\Delta t})$$

Accurate knowledge of the vectors  $\phi_{j,t-\Delta t}$  and  $\phi_{j,t-2\Delta t}$  are required to yield an accurate solution for  $\phi_{j,t+\Delta t}$  and



normally the Houbolt integration scheme requires a special starting procedure to determine the initial two vectors  $\phi_{j,\Delta t}$  and  $\phi_{j,2\Delta t}$ . However, since the primary interest of this problem is to integrate the equations until they converge to a steady state solution, it is not necessary to obtain an accurate time history of the flow. Errors induced by the inaccurate starting vectors will vanish as time approaches infinity. Therefore, the starting vectors may be chosen somewhat arbitrarily.





## VII. CONVERGING-DIVERGING NOZZLE

S. F. Shen [Ref. 10] demonstrated the feasibility of calculating compressible flows through a converging-diverging (Laval) nozzle by dividing the region of calculations into three patches, a subsonic region, a supersonic region and a transonic one, of course bounded by the other two regions. The locations of the boundaries for each region were chosen arbitrarily provided the sonic line is bracketed by the subsonic and supersonic boundaries.

Two different finite element formulations were used for the subsonic and supersonic regions, but Shen [10] resorted to analytical approximations to cover the transonic patch. This restricted the calculations to nozzles with small throat curvatures because no analytic solutions exist for nozzles with large throat curvatures. It is conceivable that STRANL-II could be adapted to provide a continuous solution throughout all three regions.

Outside the transonic region of flow the governing small disturbance equation is

$$(1 - M_{\infty}^2)\phi_{xx} + \phi_{yy} = 0 \quad (29)$$

This holds for both subsonic and supersonic flow. Comparing equation 29 with equation 1, the transonic small disturbance equation, we notice that only the non-linear coefficient  $M_{\infty}^2(\gamma + 1)u$  distinguishes the two equations from each other.



This coefficient becomes negligible when the Mach number becomes less than .8 or greater than 1.2. With this consideration in mind, it was assumed that equation 1 would adequately describe the flow through the nozzle and that the finite element formulation developed for the non-lifting airfoil would apply to the Laval nozzle.

#### A. BOUNDARY CONDITIONS

Two solutions are possible for a converging-diverging nozzle: 1) Symmetric flow, where the flow is subsonic through the domain, except for a small supersonic region near the wall in the throat, and 2) Asymmetric solution, where the flow accelerates to sonic velocity in the throat and then continues to accelerate to supersonic velocity in the diverging section. Different boundary conditions apply for the two solutions. For the symmetric case, both inlet and exit velocities must be specified. Inlet and exit velocities are equivalent in the subsonic solution. The supersonic solution requires that only the inlet velocities be specified. If the exit velocities are also applied, the problem is overspecified and the solution may not converge.

Velocities at the inlet and exit are not uniform in the y direction, therefore the disturbances cannot be set to zero as in the case of the non-lifting airfoil. Boundary velocities must be calculated by solving equation 29 analytically.

Equation 29 is a linear equation which can be mapped to Laplace's equation,

$$\nabla^2 \phi = 0,$$



by letting  $y' = \sqrt{(1 - M_\infty^2)} y$ . Laplace's equation is easily solved for the case of the hyperbolic nozzle by transforming from cartesian coordinates to elliptic coordinates. This transformation simplifies the solution because the stream lines must be hyperbolas to follow the nozzle boundary and therefore follow the hyperbolic coordinate  $v = \text{constant}$ .

If the elliptic coordinates  $\mu$  and  $v$  are chosen such that the curves  $\mu = \text{constant}$  are ellipses and the  $v$  curves are hyperbolas, then the velocity potential which satisfies Laplace's equation for a hyperbolic nozzle is simply

$$\phi = A\mu$$

where  $A$  is a constant of integration. The stream function is

$$\psi = Av$$

The transformation  $w = \mu + iv = \cosh^{-1}(2z/a)$  gives rise to the elliptic coordinates

$$y = 1/2 a \cosh \mu \cos v, \quad x = 1/2 a \sinh \mu \sin v$$

$$r = 1/2 a [\cosh^2 \mu - \sin^2 v]$$

$$r_1 = \sqrt{(y + a/2)^2 + x^2}$$

$$r_2 = \sqrt{(y - a/2)^2 + x^2}$$

Solving for  $\mu$  and  $v$  produces

$$\mu = \cosh^{-1} [(r_1 + r_2)/a]$$

$$v = \cos^{-1} [(r_1 - r_2)/a]$$

The nozzle boundary is defined by  $v_0 = \text{constant}$ , which along with the equation for the nozzle wall in cartesian



coordinates,  $y' = f(x)$ , implicitly defines the constant  $a$ . Substituting for  $\mu$  in the velocity potential produces,

$$= A \cosh^{-1} [(r_1 + r_2)/a]$$

from which the velocities may be determined.

$$u = \phi_x = \frac{A}{\sqrt{[(r_1 + r_2)/a]^2 - 1}} \left\{ a \frac{\partial r_1}{\partial x} + a \frac{\partial r_2}{\partial y} \right\}$$

$$v = \phi_y = \frac{A}{\sqrt{[(r_1 + r_2)/a]^2 - 1}} \left\{ a \frac{\partial r_1}{\partial y} + a \frac{\partial r_2}{\partial x} \right\}$$

$$\frac{\partial r_1}{\partial x} = \frac{x}{\sqrt{(y + a/2)^2 + x^2}}$$

$$\frac{\partial r_2}{\partial x} = \frac{x}{\sqrt{(y - a/2)^2 + x^2}}$$

$$\frac{\partial r_1}{\partial y} = \frac{y + a/2}{\sqrt{(y + a/2)^2 + x^2}}$$

$$\frac{\partial r_2}{\partial y} = \frac{y - a/2}{\sqrt{(y - a/2)^2 + x^2}}$$

The constant of integration  $A$  may be determined by specifying the flow rate through the nozzle, but when the inlet velocities are normalized with respect to the freestream velocity ( $U_\infty$ )  $A$  is factored out of the problem.

Velocities for compressible flow can be solved by mapping back to the physical coordinate system ( $x, y$  plane).

Other boundary conditions are universal to both problems. These are:

$$u = 0 \quad \text{on the line of symmetry}$$

$$v = (1 + u)df/dx \quad \text{at the nozzle wall}$$





$F(x)$  defines the nozzle boundary in terms of a ratio of the throat semi-height as a function of  $x$ . The throat semi-height is taken to be 1 for convenience.

Pressure ratio, sound speed, and Mach number are calculated as before by equations 5 through 8.



## VIII. DISCUSSION OF RESULTS

### A. TIME INTEGRATION TO A STEADY STATE SOLUTION

As stated before, the Houbolt method of integration is stable for all time steps. Results of the test cases bear this out with the larger time steps providing the most rapid convergence to a steady state solution. Time steps were tried from  $t = .1$  to  $t = 100$ . Time steps larger than  $t = 100$  were not attempted because as  $t$  becomes too large the influence that the damping and mass matrices have on the effective stiffness matrix becomes negligible compared to the stiffness matrix. That is:

$$SE_{ij} = S_{ij} + 2SC_{ij}/\Delta t^2 + 6SM_{ij}/\Delta t$$

as  $\Delta t \rightarrow \infty$

$$SE_{ij} \approx S_{ij}$$

The starting solutions were chosen somewhat arbitrarily.  $\phi_{j,2\Delta t}$  was chosen to satisfy the first iteration of the steady solution

$$S_{ij}\phi_{j,2\Delta t} = 0$$

when the non-linear term  $(u)$  in the coefficient

$$1 - M_{\infty}^2 - M_{\infty}^2(\gamma + 1)u$$

was set to zero.  $\phi_{j,\Delta t}$  and  $\phi_{j,0}$  were chosen as multiples of  $\phi_{j,2\Delta t}$  and respectively they were



$$\phi_{j,\Delta t} = .5\phi_{j,2\Delta t}$$

$$\phi_{j,0} = 0$$

This starting procedure proved to be superior to choosing the first three vectors closer to the converged solution. If  $\phi_{j,2\Delta t}$ ,  $\phi_{j,\Delta t}$ , and  $\phi_{j,0}$  were chosen to be the last three time steps of the previous case, the solution oscillated and converged much slower than with the starting solutions chosen as above.

The stiffness, mass, and damping matrices were recalculated after each time step, using the under-relaxation technique described above. This was necessary to utilize the special assembly procedures invoked by STRANL-II to prevent the inadmissible influence of downwind nodes from propagating upstream in the supersonic region.

For barely critical flow ( $M_\infty = .861$ ) and subsonic flow, an under-relaxation factor  $\theta = 1$  (successive approximation) resulted in convergence to a steady state solution after only three time steps. Eleven time steps were required for the supercritical solution to converge using the same relaxation factor. Reducing  $\theta$  to .5 increased the rate of convergence and the solution achieved steady state after six time steps. Figure 4 compares the steady state solution for a 6% thick circular arc airfoil at  $M_\infty = .909$ , using the same integration method, with the results obtained in Ref. [5]. Chan's results converged in 10 iterations after using the results from the barely critical flow as an initial guess to the



supercritical solution. Figure 4 is a plot of local Mach numbers at boundary nodes on the airfoil.

## B. CONVERGING-DIVERGING NOZZLE

The nozzle chosen for the test cases was the two-dimensional Oswatitsch nozzle with the boundary defined as

$$y = 1 + \sqrt{.2(x - 2.5)^2}$$

where the throat at  $x = 2.5$  has a semi-height of 1. The inlet was taken to be  $x = 0$  and the exit was at  $x = 5$ .  $M_\infty = .44$ , the inlet Mach number on the nozzle center-line was chosen to yield sonic conditions in the throat.

Two solutions were possible for this inlet condition,- the symmetric solution and the asymmetric solution; but neither solution was achieved by the finite element method. Although the solution converged for the subsonic case in three iterations, center-line Mach numbers deviated significantly from both one-dimensional theory and from Oswatitsch's approximation [Fig. 5]. When the local Mach number  $M$  exceeded the inlet Mach number by approximately .2 ( $M \geq .64$ ) the solution was invalid. Differences at the center part of the nozzle are due to an essentially incorrect free stream Mach number. Patching the solution at  $x \approx 1.5$  would improve the solution.

A second test case was run for the supersonic section of the nozzle with the inlet boundary on the sonic line. The exit boundary was left free to float. Here the solution was unstable and no meaningful results were obtained.





— Steady solution  
+ Time integration

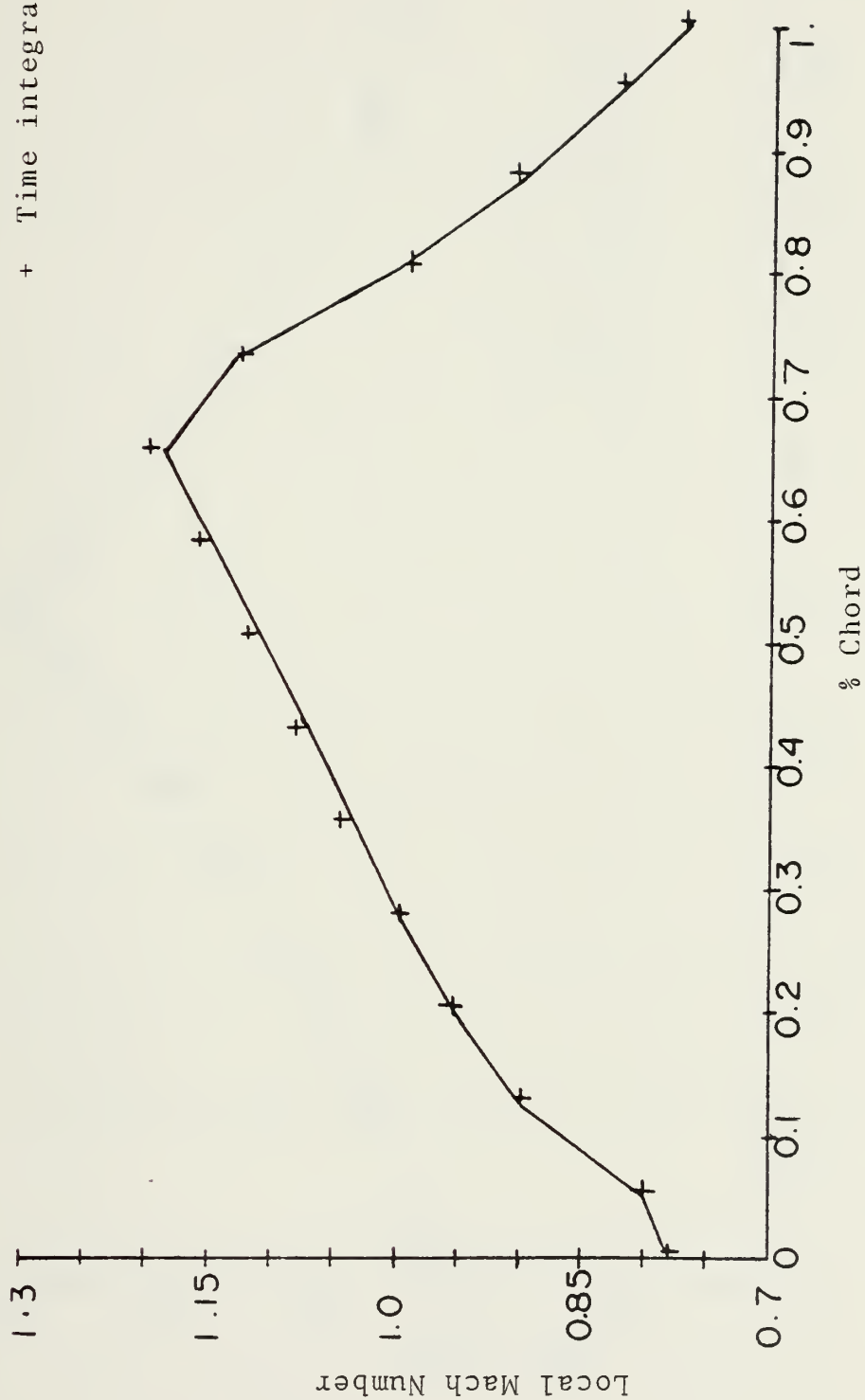


Figure 4 - Comparison of Time Integration Results with Steady State Results.



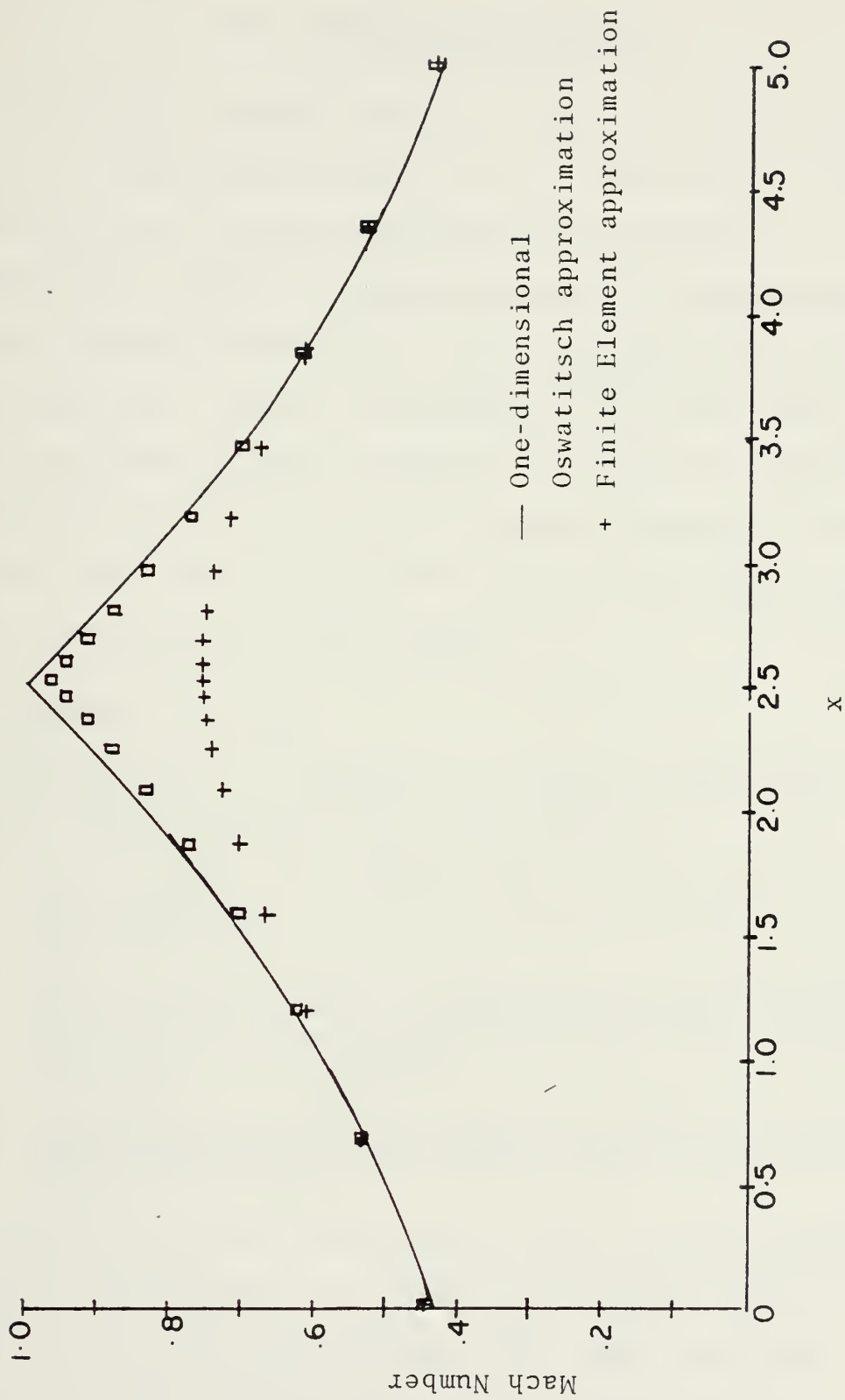


Figure 5 - Nozzle Center-Line, Inlet Mach Number 0.4300.



## IX. PROGRAM MODIFICATION

The finite element computer program for non-lifting air-voils, as developed in Ref. [5], is separated into two parts. These have been designated STRANL-I and STRANL-II by Lockheed Corporation. STRANL-I generates a finite element mesh to be used as inputs to STRANL-II, which assembles the finite element equations, applies the boundary conditions, and solves the non-linear system of equations. Detailed descriptions and instructions for the use of the two programs can be found in Ref. [5]. Only the modifications to the above programs will be discussed in this section.

### A. UNSTEADY EQUATIONS

Modifications to STRANL-II to form and solve the unsteady finite element equations were three-fold:

- 1) The new elemental matrices  $SC_{ij}$  and  $SM_{ij}$  were calculated and assembled.
- 2) All the matrices were stored on an external magnetic disk to be accessed and reassembled later because of the amount of space required to store three large matrices, in core memory.
- 3) The effective stiffness matrix  $SE_{ij}$  and the effective load vector  $RE_{ij}$  were assembled, and the system of equations solved.

Several existing subroutines in the original STRANL-II program were modified to assemble the damping and mass matrices. These include subroutines NEWK, EMTC, DERV, and EMQT. Two new subroutines were added to perform the other tasks.



## B. MODIFICATIONS TO CALCULATE THE MASS AND DAMPING MATRICES

EMTC in the STRANL-II program calculated the elemental stiffness matrix by numerically integrating the equation,

$$S_{ij} = Q_i P_j dA$$

Equations were added to EMTC to perform the additional numerical integrations for the mass and damping matrices. All three matrices were calculated at the same time. EMQT assembled the elemental stiffness matrices for a quadrilateral and trapezoidal element from the contributions of the triangular elements. Mass and damping matrices were calculated in the same fashion.

Subroutine NEWK, an original subroutine in STRANL-II, which assembled the finite element equations for steady flow, was modified to assemble  $SC_{ij}$  and  $SM_{ij}$ . A new calling argument, NMAT was passed to NEWK, which assembled contributions from the triangular, quadrilateral and trapezoidal element matrices into the global matrices  $S_{ij}$ ,  $SC_{ij}$ , and  $SM_{ij}$ , depending on NMAT being 1, 2 or 3.

### 1. Subroutine STORE

Given a non-symmetric matrix stored in a banded node, plus the right hand side vector, subroutine STORE separates this system into two matrices and stores them on a magnetic disk. Figures 6 and 7 show the decomposition of a banded matrix into banded storage, and the further decomposition of this banded stored matrix to two smaller matrices by subroutine STORE. In these figures, D, L, and U represent the





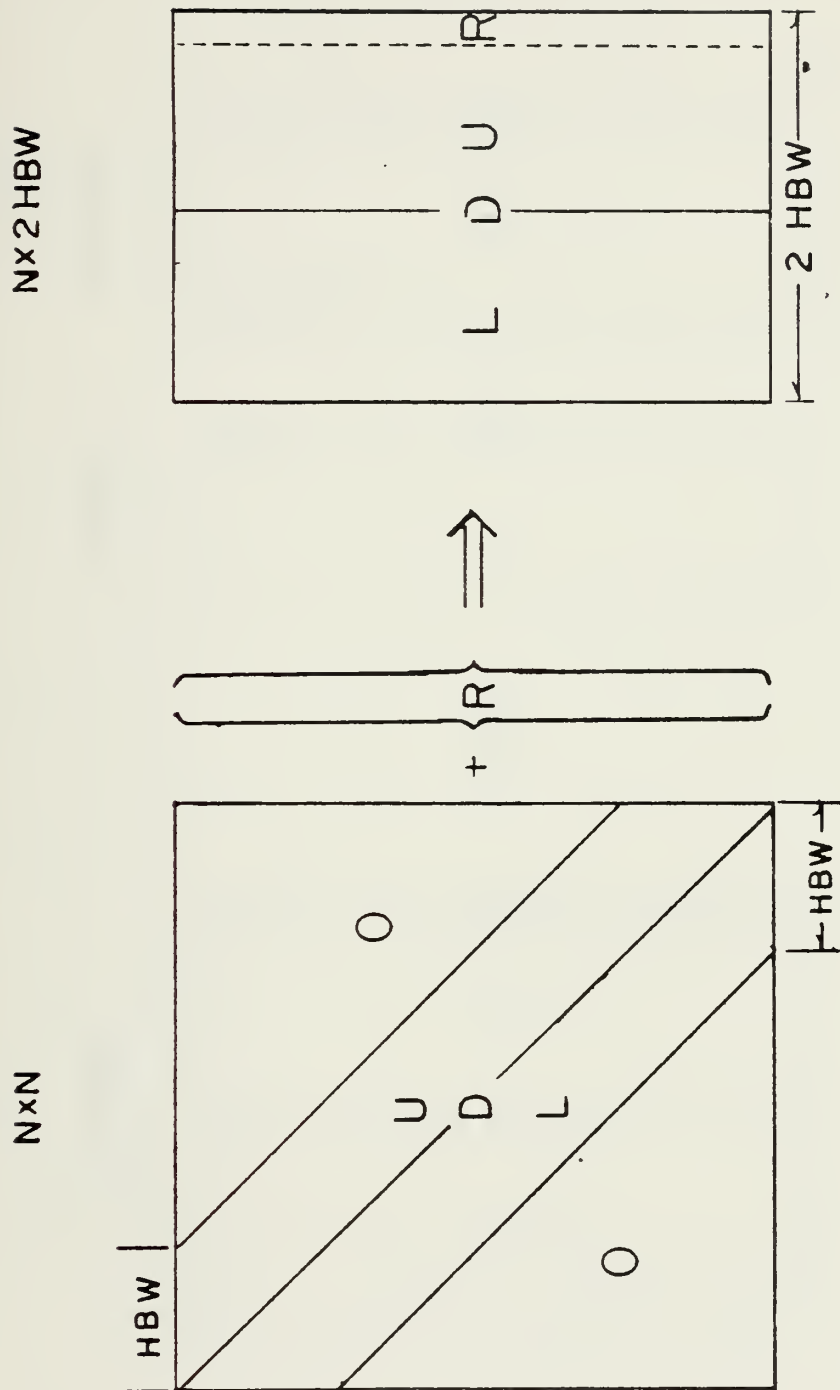


Figure 6 - Decomposition of a Banded Matrix Plus Right Hand Side



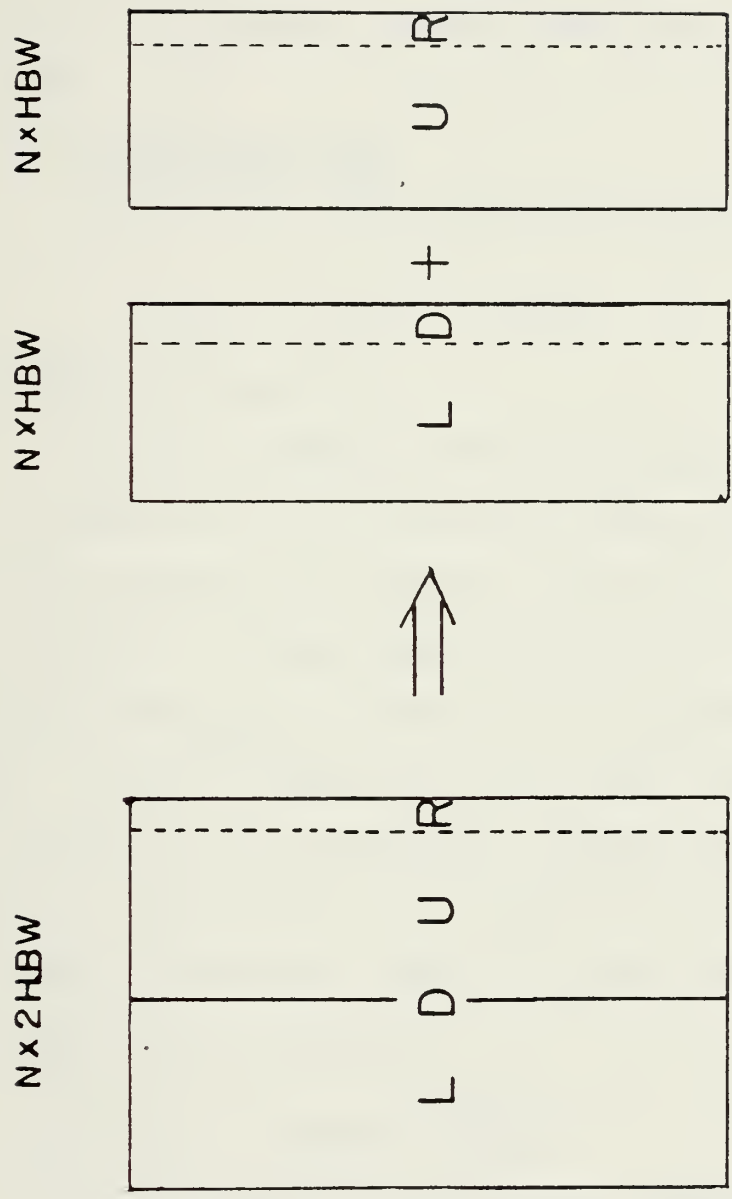


Figure 7 - Separation of a Banded Stored Matrix by Store



diagonal matrix, the lower triangular matrix, and the upper triangular matrix respectively. HBW is the half bandwidth and R is the right hand side vector.

STORE requires an additional work area one-half the size of the originally dimensioned matrix which is to be stored.

## 2. Subroutine TIME

Subroutine TIME integrates the system

$$S_{ij}\phi_j + SC_{ij}\dot{\phi}_j + SM_{ij}\ddot{\phi}_j = R_j$$

by Houbolt integration.

TIME reassembles the three matrices which were stored on the magnetic disk to form the effective stiffness matrix and the effective load vector. Once this system of equations is assembled, a banded equation solver is called to yield the solution for  $\phi_{j,t+\Delta t}$ . Figure 8 is a schematic flow chart of TIME. In Fig. 8 when  $L = 1$ , the lower triangular matrix and the diagonal of the effective stiffness matrix are formed by adding the appropriate contributions from the stiffness, mass and damping matrices. When  $L = 2$ , the upper triangular matrix is formed in like fashion.

## C. CONVERGING-DIVERGING NOZZLE

### 1. Application of the Boundary Conditions

Regardless of the type of problem for which a set of system equations have been assembled, the equations will have the form

$$K_{ij} x_i = R_i$$



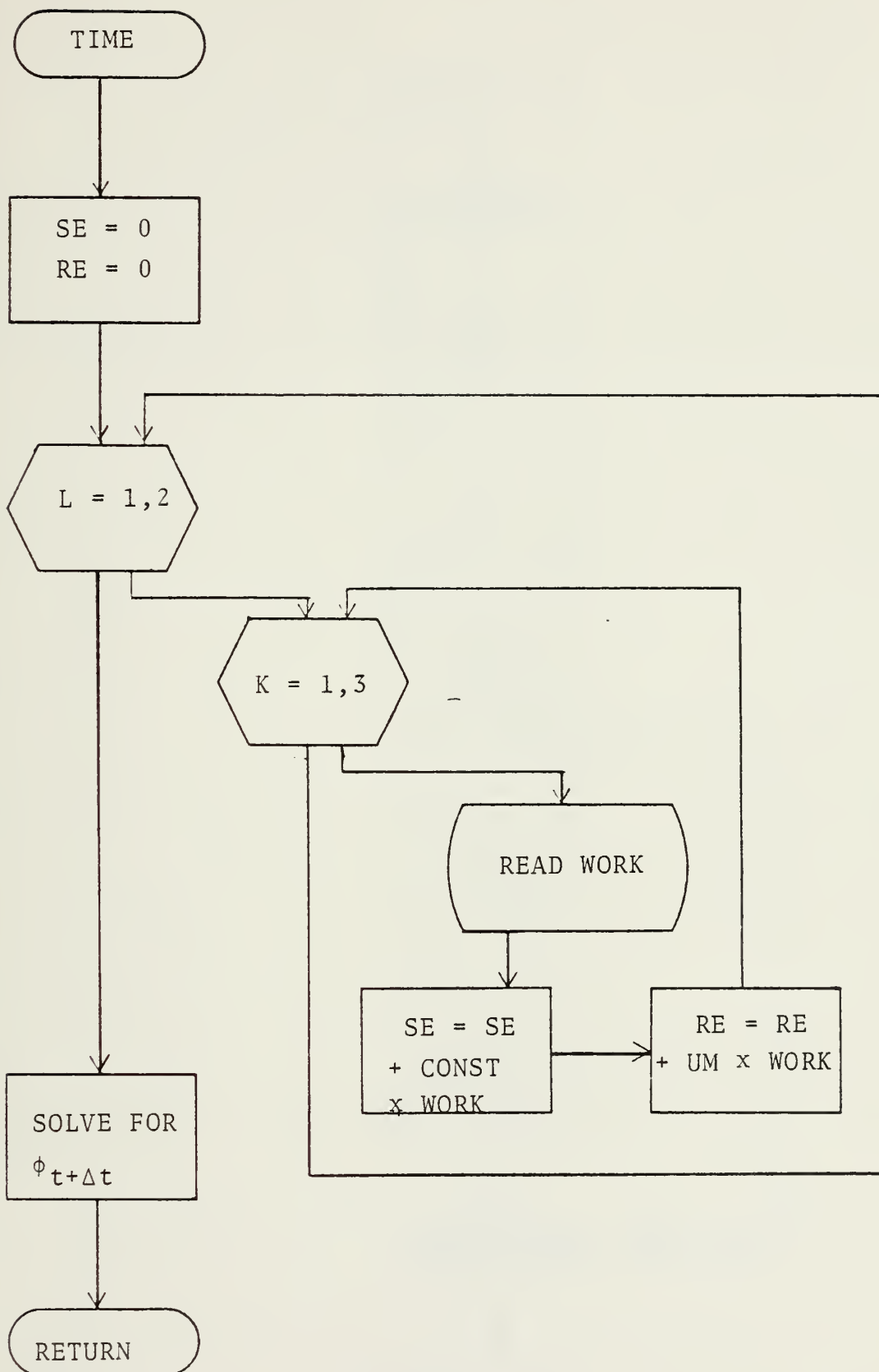


Figure 8 - Flow Chart of Subroutine TIME





STRANL-II

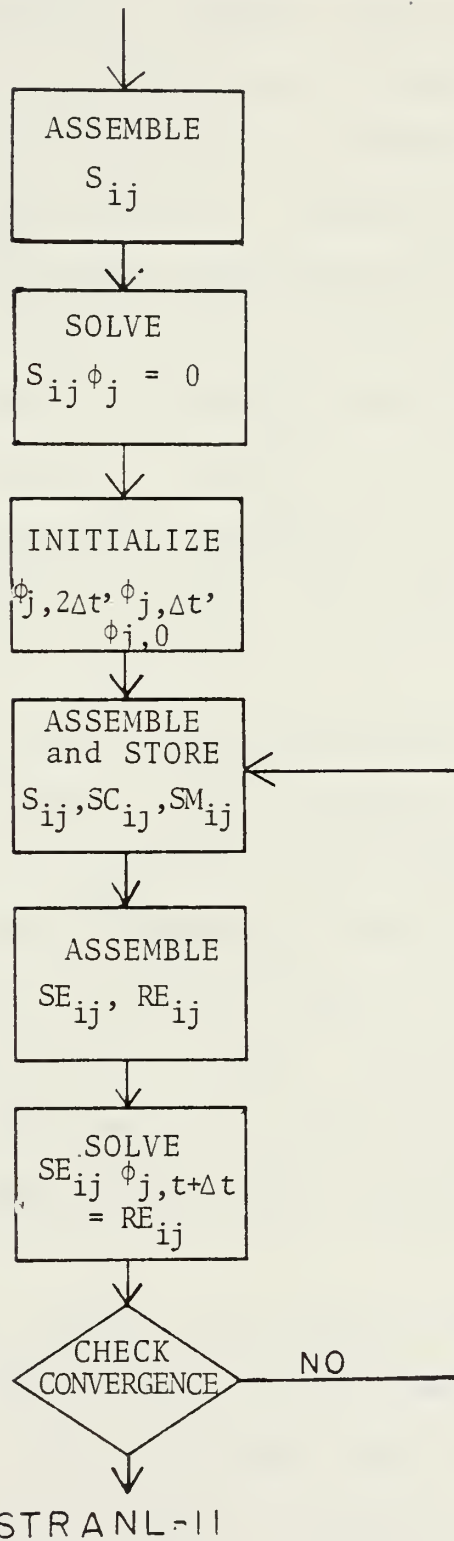


Figure 9 - Flow Chart of STANL-II Modification to Integrate Unsteady Equations



in which  $K_{ij}$  is an  $n \times n$  matrix and  $x_i$  and  $R_i$  are vectors of length  $n$ . These equations do not take into account the known values of  $x_i$  on the boundaries. However, for a unique solution of the above equation, at least one or more nodal variables must be specified and  $K_{ij}$  must be modified to render it non-singular. For each equation  $i$ , either  $x_i$  or  $R_i$  must be specified but it is physically impossible to specify both  $x_i$  and  $R_i$ . There are a number of ways to apply the boundary conditions to the equations and when they are applied the number of equations is reduced. However, it is convenient to leave the number of equations unchanged to avoid major restructuring of the computer storage. One such method is described below.

If  $k$  is the subscript of the prescribed nodal variable, the  $k^{\text{th}}$  row and the  $k^{\text{th}}$  column of the original  $K_{ij}$  matrix are set to zero,  $K_{kk}$  is set to 1 and  $R_k$  is replaced by the known value of  $x_k$ . Each of the  $n-1$  remaining terms of  $R_i$  is modified by subtracting from it the value  $K_{ik}x_k$ . This procedure is repeated for all the boundary values. Of course, when the matrix is stored in a banded mode, the algorithm will differ from that for the  $n \times n$  square matrix, but the procedure is similar.

Subroutine BNDRY applies the boundary conditions for the modified program. Setting the option parameter IOPT(4), in STRANL-II, equal to 1 will call BNDRY which will read the boundary velocities and apply them to a banded stored matrix by the method described above.



When the value of the  $x_k$  is zero, as in the non-lifting airfoil problem, the algorithm becomes simpler than the above method because there is no need to either set the  $k^{\text{th}}$  column to zero or subtract the value  $K_{ik}x_k$  from the right hand vector.



## X. TEST CASES

In computing the flow field for either the non-lifting airfoil or for the Laval nozzle, the following procedures were followed:

- 1) The desired mesh was sketched with each node assigned a number.
- 2) Appropriate input cards based on the sketch were prepared and supplied to STRANL-I to generate the data on punched cards for input to STRANL-II.
- 3) The above punched cards were supplied to STRANL-II with three additional cards as input parameters for each case, plus additional cards for the boundary velocities, if the nozzle solution is desired.
- 4) Results of the finite element calculations are printed after each iteration, and the converged solution is punched on cards for possible later use.

### A. TIME INTEGRATION TO A STEADY STATE SOLUTION

Test cases for the integration of the unsteady transonic finite element equations were conducted to calculate the steady transonic flow over a 6% thick circular arc airfoil. These tests were made using the same airfoil, mesh, and free-stream Mach numbers as Chan et al. [Ref. 5] published. These conditions were chosen to provide a source for comparison of the results.

Freestream Mach numbers used in these calculations were:

$$M_{\infty} = .806 \quad (\text{subcritical})$$

$$M_{\infty} = .861 \quad (\text{barely critical})$$

$$M_{\infty} = .909 \quad (\text{supercritical}).$$





Each case was treated individually with  $\phi_j = 0$  used as the initial guess for each case, whereas Chan et al. [Ref. 5] used zero for the initial guess for  $M_\infty = .806$  and then used the computed results as the initial guess for each subsequent case.

DELTA, the value for the time step, is input by a parameter specified in columns 41-45 of the second card following the title card for each case when the unsteady option (IOPT(6) = 1) is selected.

### 1. STRANL-I Program

#### a. Input

Input cards used to generate the finite element mesh are listed on the next page. Cards were arranged in accordance with Ref. [5], in the following order:

- Title card
- Option card
- Element cards
- Card for the total number of nodes
- Node coordinate cards
- Card for the number of boundary nodes
- Card for the boundary nodes at infinity
- Card for the nodes on the line of symmetry
- Card for the nodes on the airfoil
- Cards for the slope of the airfoil.

Input cards to STRANL-I for these tests are listed on the next three pages.

#### b. Output

Output from STRANL-I is in the form of printouts and punched cards. Printouts from STRANL-I are listed on the following eight pages.



# INPUT TO STRANL-I

STEADY	TRANSCNIC	FLOW	MESH	6--154	ELEMENTS,	170	NODES
3	4	1	0	21	26	7	TRA00010
3	4	1	0	27	32	32	TRA00020
3	4	1	0	28	33	32	TRA00030
4	3	1	0	33	36	36	TRA00040
3	4	1	0	34	37	36	TRA00050
4	3	1	0	49	53	52	TRA00060
4	3	1	0	52	55	56	TRA00070
4	3	1	0	56	57	56	TRA00080
4	3	1	0	62	60	61	TRA00090
3	4	1	0	65	62	66	TRA00100
4	3	1	0	66	67	66	TRA00110
4	3	1	0	62	67	67	TRA00120
4	3	1	0	65	67	88	TRA00130
4	3	1	0	62	67	88	TRA00140
4	3	1	0	65	67	92	TRA00150
4	3	1	0	62	67	97	TRA00160
4	3	1	0	65	67	103	TRA00170
4	3	1	0	62	67	110	TRA00180
4	3	1	0	65	67	117	TRA00190
4	3	1	0	62	67	124	TRA00200
4	3	1	0	65	67	131	TRA00210
4	3	1	0	62	67	147	TRA00220
4	3	1	0	65	67	150	TRA00230
4	3	1	0	62	67	154	TRA00240
4	3	1	0	65	67	154	TRA00250
4	3	1	0	62	67	159	TRA00260
4	3	1	0	65	67	170	TRA00270
4	3	1	0	62	67	175	TRA00280
4	3	1	0	65	67	180	TRA00290
4	3	1	0	62	67	194	TRA00300
4	3	1	0	65	67	199	TRA00310
4	3	1	0	62	67	205	TRA00320
4	3	1	0	65	67	210	TRA00330
4	3	1	0	62	67	215	TRA00340
4	3	1	0	65	67	223	TRA00350
4	3	1	0	62	67	237	TRA00360
4	3	1	0	65	67	249	TRA00370
4	3	1	0	62	67	258	TRA00380
4	3	1	0	65	67	263	TRA00390
4	3	1	0	62	67	268	TRA00400
4	3	1	0	65	67	270	TRA00410
4	3	1	0	62	67	275	TRA00420
4	3	1	0	65	67	280	TRA00430
4	3	1	0	62	67	286	TRA00440
4	3	1	0	65	67	293	TRA00450
4	3	1	0	62	67	298	TRA00460
4	3	1	0	65	67	304	TRA00470
4	3	1	0	62	67	309	TRA00480
4	3	1	0	65	67	315	TRA00490
4	3	1	0	62	67	318	TRA00500
4	3	1	0	65	67	323	TRA00510
4	3	1	0	62	67	328	TRA00520
4	3	1	0	65	67	333	TRA00530
4	3	1	0	62	67	338	TRA00540
4	3	1	0	65	67	343	TRA00550
4	3	1	0	62	67	348	TRA00560
4	3	1	0	65	67	353	TRA00570
4	3	1	0	62	67	358	TRA00580
4	3	1	0	65	67	363	TRA00590
4	3	1	0	62	67	368	TRA00600
4	3	1	0	65	67	373	TRA00610
4	3	1	0	62	67	378	TRA00620
4	3	1	0	65	67	383	TRA00630
4	3	1	0	62	67	388	TRA00640
4	3	1	0	65	67	393	TRA006







86	89	93	98	104	111	118	125	132	135	146	153	159	164	168
.06022	+0.10820	+0.09004	+0.09004	+0.09004	+0.07193	+0.07193	+0.05385	+0.05385	+0.03589	+0.03589	+0.01794	+0.01794	+0.00000	+0.00000
-.01794	-0.03589	-0.05388	-0.05388	-0.05388	-0.07193	-0.07193	-0.09004	-0.09004	-0.10820	-0.10820	-0.01794	-0.01794	-0.00000	-0.00000
1.0	8.0	0.1	0.1	0.1	3.5	3.5	0.1	0.1	0.1	0.1	0.1	0.1	0.1	0.1

CIRCULAR ARC









[illegible]





















120	01	0	2438E	0	0	00
121	01	0	58515E	0	0	00
122	01	0	55406E	0	0	00
123	01	0	75676E	0	0	01
124	01	0	10000E	0	0	01
125	01	0	30000E	0	0	00
126	01	0	12570E	0	0	00
127	01	0	24534E	0	0	00
128	01	0	39290E	0	0	00
129	01	0	56518E	0	0	00
130	01	0	77192E	0	0	00
131	01	0	10200E	0	0	01
132	01	0	29330E	0	0	01
133	01	0	12768E	0	0	00
134	01	0	24438E	0	0	00
135	01	0	38515E	0	0	00
136	01	0	55406E	0	0	00
137	01	0	75676E	0	0	00
138	01	0	10000E	0	0	01
139	01	0	27310E	0	0	01
140	01	0	12124E	0	0	00
141	01	0	23395E	0	0	00
142	01	0	36521E	0	0	00
143	01	0	53151E	0	0	00
144	01	0	72628E	0	0	00
145	01	0	96000E	0	0	00
146	01	0	23940E	0	0	01
147	01	0	11418E	0	0	00
148	01	0	22246E	0	0	00
149	01	0	35241E	0	0	00
150	01	0	50834E	0	0	00
151	01	0	65546E	0	0	00
152	01	0	92000E	0	0	00
153	01	0	19220E	0	0	01
154	01	0	11742E	0	0	00
155	01	0	23526E	0	0	00
156	01	0	37668E	0	0	00
157	01	0	54637E	0	0	00
158	01	0	75000E	0	0	00
159	01	0	13150E	0	0	01
160	01	0	11316E	0	0	00
161	01	0	23317E	0	0	00
162	01	0	37718E	0	0	00
163	01	0	55000E	0	0	00
164	01	0	72200E	0	0	02
165	01	0	10030E	0	0	00
166	01	0	21380E	0	0	00
167	01	0	35000E	0	0	00



168	15	0.30000E 01	0.0						
169	14	0.30000E 01	0.10000E 00						
170	13	0.30000E 01	0.22000E 00						
NODES AT FARFIELD									
1	6	11	16	21	26	31	32	36	39
2	66	71	76	81					
NODES CN LINE OF SYMMETRY									
3	5	82	83	84	85				
NODES CN THE AIRFOIL									
4	89	93	98	104	111	118	125	132	139
NODES ON THE AIRFOIL									
SLCFE ALONG NODES									
5	6220E-01	01	0.10820E 00	0.90040E-01	0.0	0.71930E-01	0.71930E-01	0.53890E-01	0.53890E-01
6	35890E-01	01	0.17940E-01	0.0	0.0	0.0	0.0	0.0	0.0
7	17940E-01	01	-0.35890E-01	-0.53880E-01	-0.53880E-01	-0.53880E-01	-0.53880E-01	-0.53880E-01	-0.53880E-01
8	10820E 00	00	-0.60220E-01	-0.60220E-01	-0.60220E-01	-0.60220E-01	-0.60220E-01	-0.60220E-01	-0.60220E-01
-C.10820E 00									

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## 2. STRANL-II Program

### a. Input

Listed on the next three pages are the input cards to the STANL-II program.

### b. Output

The output from this program is in the form of printouts for each iteration and punched cards for the converged solution. Output from STRANL-II for  $M_\infty = .909$  is listed on the following eight pages.



# INPUT TO STRANL-II

TIME	INTEGRATION	SOLUTION	--	CIRCULAR	ARC	--	170	NODES	--	M=0.909
0 0	154	170	0.909	8	15	500	1	1	1	1
1 3	11	17	22	18	14	10	33	38	44	50
1 1	17	13	27	33	38	44	50	53	57	63
1 5	22	33	44	50	53	57	63	68	73	78
1 9	27	33	44	50	53	57	63	68	73	78
2 2	33	44	50	53	57	63	68	73	78	83
2 6	38	44	50	53	57	63	68	73	78	83
2 9	44	50	53	57	63	68	73	78	83	88
3 3	50	53	57	63	68	73	78	83	88	93
3 7	55	63	68	73	78	83	88	93	98	103
3 9	63	68	73	78	83	88	93	98	103	108
4 0	68	73	78	83	88	93	98	103	108	113
4 4	73	78	83	88	93	98	103	108	113	118
4 8	78	83	88	93	98	103	108	113	118	123
5 1	83	88	93	98	103	108	113	118	123	128
5 5	88	93	98	103	108	113	118	123	128	133
5 6	93	98	103	108	113	118	123	128	133	138
6 2	98	103	108	113	118	123	128	133	138	143
6 7	103	108	113	118	123	128	133	138	143	148
7 2	108	113	118	123	128	133	138	143	148	153
7 7	113	118	123	128	133	138	143	148	153	158
7 8	118	123	128	133	138	143	148	153	158	163
8 3	123	128	133	138	143	148	153	158	163	168
8 7	128	133	138	143	148	153	158	163	168	173
9 3	133	138	143	148	153	158	163	168	173	178
9 7	138	143	148	153	158	163	168	173	178	183
1 0	143	148	153	158	163	168	173	178	183	188
1 3	148	153	158	163	168	173	178	183	188	193
1 7	153	158	163	168	173	178	183	188	193	198
2 0	158	163	168	173	178	183	188	193	198	203
2 4	163	168	173	178	183	188	193	198	203	208
2 8	168	173	178	183	188	193	198	203	208	213
3 2	173	178	183	188	193	198	203	208	213	218
3 6	178	183	188	193	198	203	208	213	218	223
4 0	183	188	193	198	203	208	213	218	223	228
4 4	188	193	198	203	208	213	218	223	228	233
4 8	193	198	203	208	213	218	223	228	233	238
5 2	198	203	208	213	218	223	228	233	238	243
5 6	203	208	213	218	223	228	233	238	243	248
6 0	208	213	218	223	228	233	238	243	248	253
6 4	213	218	223	228	233	238	243	248	253	258
6 8	218	223	228	233	238	243	248	253	258	263
7 2	223	228	233	238	243	248	253	258	263	268
7 6	228	233	238	243	248	253	258	263	268	273
8 0	233	238	243	248	253	258	263	268	273	278
8 4	238	243	248	253	258	263	268	273	278	283
8 8	243	248	253	258	263	268	273	278	283	288
9 2	248	253	258	263	268	273	278	283	288	293
9 6	253	258	263	268	273	278	283	288	293	298
1 0 0	258	263	268	273	278	283	288	293	298	303
1 0 4	263	268	273	278	283	288	293	298	303	308
1 0 8	268	273	278	283	288	293	298	303	308	313
1 1 2	273	278	283	288	293	298	303	308	313	318
1 1 6	278	283	288	293	298	303	308	313	318	323
1 2 0	283	288	293	298	303	308	313	318	323	328
1 2 4	288	293	298	303	308	313	318	323	328	333
1 2 8	293	298	303	308	313	318	323	328	333	338
1 3 2	298	303	308	313	318	323	328	333	338	343
1 3 6	303	308	313	318	323	328	333	338	343	348
1 4 0	308	313	318	323	328	333	338	343	348	353
1 4 4	313	318	323	328	333	338	343	348	353	358
1 4 8	318	323	328	333	338	343	348	353	358	363
1 5 2	323	328	333	338	343	348	353	358	363	368
1 5 6	328	333	338	343	348	353	358	363	368	373
1 6 0	333	338	343	348	353	358	363	368	373	378
1 6 4	338	343	348	353	358	363	368	373	378	383
1 6 8	343	348	353	358	363	368	373	378	383	388
1 7 2	348	353	358	363	368	373	378	383	388	393
1 7 6	353	358	363	368	373	378	383	388	393	398
1 8 0	358	363	368	373	378	383	388	393	398	403
1 8 4	363	368	373	378	383	388	393	398	403	408
1 8 8	368	373	378	383	388	393	398	403	408	413
1 9 2	373	378	383	388	393	398	403	408	413	418
1 9 6	378	383	388	393	398	403	408	413	418	423
2 0 0	383	388	393	398	403	408	413	418	423	428
2 0 4	388	393	398	403	408	413	418	423	428	433
2 0 8	393	398	403	408	413	418	423	428	433	438
2 1 2	398	403	408	413	418	423	428	433	438	443
2 1 6	403	408	413	418	423	428	433	438	443	448
2 2 0	408	413	418	423	428	433	438	443	448	453
2 2 4	413	418	423	428	433	438	443	448	453	458
2 2 8	418	423	428	433	438	443	448	453	458	463
2 3 2	423	428	433	438	443	448	453	458	463	468
2 3 6	428	433	438	443	448	453	458	463	468	473
2 4 0	433	438	443	448	453	458	463	468	473	478
2 4 4	438	443	448	453	458	463	468	473	478	483
2 4 8	443	448	453	458	463	468	473	478	483	488
2 5 2	448	453	458	463	468	473	478	483	488	493
2 5 6	453	458	463	468	473	478	483	488	493	498
2 6 0	458	463	468	473	478	483	488	493	498	503
2 6 4	463	468	473	478	483	488	493	498	503	508
2 6 8	468	473	478	483	488	493	498	503	508	513
2 7 2	473	478	483	488	493	498	503	508	513	518
2 7 6	478	483	488	493	498	503	508	513	518	523
2 8 0	483	488	493	498	503	508	513	518	523	528
2 8 4	488	493	498	503	508	513	518	523	528	533
2 8 8	493	498	503	508	513	518	523	528	533	538
2 9 2	498	503	508	513	518	523	528	533	538	543
2 9 6	503	508	513	518	523	528	533	538	543	548
3 0 0	508	513	518	523	528	533	538	543	548	553
3 0 4	513	518	523	528	533	538	543	548	553	558
3 0 8	518	523	528	533	538	543	548	553	558	563
3 1 2	523	528	533	538	543	548	553	558	563	568
3 1 6	528	533	538	543	548	553	558	563	568	573
3 2 0	533	538	543	548	553	558	563	568	573	578
3 2 4	538	543	548	553	558	563	568	573	578	583
3 2 8	543	548	553	558	563	568	573	578	583	588
3 3 2	548	553	558	563	568	573	578	583	588	593
3 3 6	553	558	563	568	573	578	583	588	593	598
3 4 0	558	563	568	573	578	583	588	593	598	603
3 4 4	563	568	573	578	583	588	593	598	603	608
3 4 8	568	573	578	583	588	593	598	603	608	613
3 5 2	573	578	583	588	593	598	603	608	613	618
3 5 6	578	583	588	593	598	603	608	613	618	623
3 6 0	583	588	593	598	603	608	613	618	623	628
3 6 4										







111  
161 3150  
14097  
187 6536  
133  
173 303  
14108  
11866  
137  
126  
172 292  
14259  
11877  
11844  
125  
171 289  
14360  
11508  
11645  
124  
183 3920  
14974  
11969  
11946  
123  
182 432  
15018  
119570  
11947  
122  
181 420  
15156  
11974  
11983  
138  
193 418  
15207  
11975  
11949  
11914  
137  
192 407  
15818  
11987  
11964  
11915  
136  
191 3818  
15977  
11997  
11975  
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11924  
135  
103 3323  
16028  
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11785  
11965  
11925  
134  
102 2216  
12949  
11047  
11795  
11972  
11926  
148  
101 5119  
13050  
11058  
11080  
11958  
11927  
147  
114 5017  
13164  
11068  
11084  
11959  
11933  
146  
113 336  
11127  
11075  
11085  
11960  
11934  
145  
112 2225  
13508  
11086  
11964  
11955  
11935





TIME INTEGRATION TO STEADY SOLUTION -- CIRCULAR ARC -- 170 NODES --  $M=0.909$   
 CONVERGENCE LIMIT = 0.0050

C C C C C 1 C C C C C C C C C C C C C C C C

NO. OF ELEMENTS = 154 NO. OF NODES = 170 FULL BANDWIDTH = 84

ELE. NO. ARC ELEMENT NCCES

1	31	21	26	C
2	11	11	11	11
3	11	11	11	11
4	11	11	11	11
5	11	11	11	11
6	11	11	11	11
7	11	11	11	11
8	11	11	11	11
9	11	11	11	11
10	11	11	11	11
11	11	11	11	11
12	11	11	11	11
13	11	11	11	11
14	11	11	11	11
15	11	11	11	11
16	11	11	11	11
17	11	11	11	11
18	11	11	11	11
19	11	11	11	11
20	11	11	11	11
21	11	11	11	11
22	11	11	11	11
23	11	11	11	11
24	11	11	11	11
25	11	11	11	11
26	11	11	11	11
27	11	11	11	11
28	11	11	11	11
29	11	11	11	11
30	11	11	11	11
31	11	11	11	11
32	11	11	11	11
33	11	11	11	11
34	11	11	11	11
35	11	11	11	11
36	11	11	11	11
37	11	11	11	11
38	11	11	11	11
39	11	11	11	11
40	11	11	11	11
41	11	11	11	11
42	11	11	11	11
43	11	11	11	11
44	11	11	11	11
45	11	11	11	11
46	11	11	11	11
47	11	11	11	11
48	11	11	11	11
49	11	11	11	11
50	11	11	11	11
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52	11	11	11	11
53	11	11	11	11
54	11	11	11	11
55	11	11	11	11
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57	11	11	11	11
58	11	11	11	11
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61	11	11	11	11
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63	11	11	11	11
64	11	11	11	11
65	11	11	11	11
66	11	11	11	11
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71	11	11	11	11
72	11	11	11	11
73	11	11	11	11
74	11	11	11	11
75	11	11	11	11
76	11	11	11	11
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81	11	11	11	11
82	11	11	11	11
83	11	11	11	11
84	11	11	11	11
85	11	11	11	11
86	11	11	11	11
87	11	11	11	11
88	11	11	11	11
89	11	11	11	11
90	11	11	11	11
91	11	11	11	11
92	11	11	11	11
93	11	11	11	11
94	11	11	11	11
95	11	11	11	11
96	11	11	11	11
97	11	11	11	11
98	11	11	11	11
99	11	11	11	11
100	11	11	11	11
101	11	11	11	11
102	11	11	11	11
103	11	11	11	11
104	11	11	11	11
105	11	11	11	11
106	11	11	11	11
107	11	11	11	11
108	11	11	11	11
109	11	11	11	11
110	11	11	11	11
111	11	11	11	11
112	11	11	11	11
113	11	11	11	11
114	11	11	11	11
115	11	11	11	11
116	11	11	11	11
117	11	11	11	11
118	11	11	11	11
119	11	11	11	11
120	11	11	11	11
121	11	11	11	11
122	11	11	11	11
123	11	11	11	11
124	11	11	11	11
125	11	11	11	11
126	11	11	11	11
127	11	11	11	11
128	11	11	11	11
129	11	11	11	11
130	11	11	11	11
131	11	11	11	11
132	11	11	11	11
133	11	11	11	11
134	11	11	11	11
135	11	11	11	11
136	11	11	11	11
137	11	11	11	11
138	11	11	11	11
139	11	11	11	11
140	11	11	11	11
141	11	11	11	11
142	11	11	11	11
143	11	11	11	11
144	11	11	11	11
145	11	11	11	11
146	11	11	11	11
147	11	11	11	11
148	11	11	11	11
149	11	11	11	11
150	11	11	11	11
151	11	11	11	11
152	11	11	11	11
153	11	11	11	11
154	11	11	11	11











139	70	0.26500E	01	0.27310E	01
140	69	0.26500E	01	0.12120E	00
141	68	0.26500E	01	0.23400E	00
142	67	0.26500E	01	0.36920E	00
143	66	0.26500E	C1	0.33150E	00
144	65	0.26500E	01	0.72630E	00
145	64	0.26500E	01	0.96000E	00
146	60	0.27250E	C1	0.23940E	01
147	59	0.27250E	01	0.11420E	00
148	58	0.27250E	C1	0.22250E	00
149	57	0.27250E	C1	0.35240E	00
150	56	0.27250E	01	0.50830E	00
151	55	0.27250E	C1	0.69550E	00
152	54	0.27250E	C1	0.92000E	00
153	49	0.28000E	C1	0.19220E	01
154	48	0.28000E	C1	0.11740E	00
155	47	0.28000E	01	0.23530E	00
156	46	0.28000E	01	0.37670E	00
157	45	0.28000E	01	0.54640E	00
158	44	0.28000E	01	0.75000E	00
159	37	0.28750E	C1	0.13150E	01
160	36	0.28750E	C1	0.11320E	00
161	35	0.28750E	C1	0.23320E	00
162	34	0.28750E	C1	0.37720E	00
163	33	0.28750E	01	0.55000E	00
164	27	0.29500E	01	0.57230E	02
165	26	0.29500E	C1	0.10030E	00
166	25	0.29500E	01	0.21380E	CG
167	24	0.29500E	01	0.35000E	00
168	15	0.30000E	01	0.0	00
169	14	0.30000E	C1	0.10000E	00
170	13	0.30000E	C1	0.22000E	00









$\text{NACH NUMBER} = 0.569 \quad \text{RELAX. FACTOR} = 0.5000$

AC. ITERATIONS = 7

NCLE	PHI I	UCUM	UCUM	CLT	LMAC	P/P/C	CP	DELM
1	1.704E-03	1.000E-00	1.905E-06	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
2	1.414E-04	1.001E-00	1.469E-06	8.796E-01	5.024E-01	5.842E-01	1.192E-05	0.181E-03
3	1.313E-04	1.540E-01	1.901E-05	7.147E-01	5.024E-01	5.842E-01	1.192E-05	0.332E-03
4	1.359E-05	9.243E-01	1.543E-04	7.761E-01	5.326E-01	5.383E-01	1.117E-02	1.543E-03
5	1.309E-05	1.000E-00	1.727E-04	8.284E-01	5.101E-01	5.850E-01	1.211E-03	0.0
6	1.434E-05	1.001E-00	1.571E-04	8.284E-01	5.101E-01	5.850E-01	1.211E-03	0.303E-03
7	1.402E-05	9.815E-01	1.547E-04	7.857E-01	5.033E-01	5.850E-01	1.063E-02	1.206E-04
8	1.770E-05	9.287E-01	1.287E-04	7.857E-01	5.033E-01	5.850E-01	1.063E-02	1.620E-04
9	1.200E-05	1.000E-00	1.044E-04	8.650E-01	5.090E-01	5.850E-01	1.154E-03	1.574E-03
10	1.784E-05	1.000E-00	1.337E-04	8.284E-01	5.053E-01	5.850E-01	1.154E-03	0.0
11	1.304E-05	9.963E-01	8.874E-04	8.190E-01	5.053E-01	5.878E-01	1.354E-03	0.0
12	1.604E-05	9.432E-01	2.012E-03	7.939E-01	5.020E-01	5.733E-01	1.123E-02	1.595E-03
13	1.503E-05	9.432E-01	1.862E-02	7.147E-01	5.020E-01	5.733E-01	1.123E-02	0.0
14	1.602E-05	1.000E-00	1.000E-04	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
15	1.572E-05	1.000E-00	1.532E-04	8.270E-01	5.055E-01	5.850E-01	1.257E-06	0.0
16	1.611E-05	9.743E-01	4.336E-04	8.192E-01	5.055E-01	5.878E-01	1.257E-06	0.084E-04
17	1.244E-05	9.562E-01	1.004E-04	7.857E-01	5.055E-01	5.878E-01	1.257E-06	1.814E-04
18	1.611E-05	1.000E-00	1.328E-02	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
19	1.244E-05	1.000E-00	2.350E-02	8.264E-01	5.090E-01	5.850E-01	1.257E-06	1.592E-03
20	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
21	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
22	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
23	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
24	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
25	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
26	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
27	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
28	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
29	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
30	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
31	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
32	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
33	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
34	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
35	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
36	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
37	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
38	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
39	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
40	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
41	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
42	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
43	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
44	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
45	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
46	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
47	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
48	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
49	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
50	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
51	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
52	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
53	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
54	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
55	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
56	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
57	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
58	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
59	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0
60	1.318E-05	9.587E-01	1.000E-00	8.264E-01	5.090E-01	5.850E-01	1.257E-06	0.0









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## B. CONVERGING-DIVERGING NOZZLE

Two test cases were run for the converging-diverging nozzle. The first case was for symmetric flow designed to yield sonic conditions in the throat of the nozzle. The second case dealt with supersonic flow in the diverging section by starting with the sonic line as the boundary of the nozzle mesh. Oswatitsch's two-dimensional nozzle [Ref. 9],  $y = 1 + \sqrt{2(x - 2.5)^2}$  with a semi-throat height of 1 at  $x = 2.5$  was used in both cases. Boundaries for the subsonic nozzle were at  $x = 0$  and  $x = 5$ .

### 1. Symmetric Solution

This problem was analyzed using the mesh shown in Fig. 10, which consists of 126 elements and 152 nodes. In the second card (the option card) IOPT(4) = 1 indicates that non-zero boundary velocities at the inlet and the exit will be read and applied by subroutine BNDRY. This option requires that the number of inlet and exit boundary nodes be specified in columns 36-40 of the next card. Perturbation velocities at the boundary nodes follow on the subsequent four cards. Subroutine BNDRY reads  $u$  and  $v$  respectively for the first boundary node and then continues reading  $u$  and  $v$  for each inlet and then each exit node in the order specified on the appropriate card.

### 2. Supersonic Case--Diverging Section

The mesh used for this case is sketched in Fig. 11 and input cards to STRANL-II follow on the next page. For this



case the options in effect are IOPT(4) = 1 and IOPT(5) = 1 which cause non-zero boundary velocities to be applied to the sonic line only.



## STEADY TRANSONIC FLOW CONVERGING DIVERGING NOZZLE-3 M=.375

[illegible]









113	114	115	116	117	118	119	120	121	122	123	124	125	126	127	128
129	130	131	132	133	134	135	136	137	138	139	140	141	142	143	144
145	146	147	148	149	150	151	152								
STEADY FLOW-CONVERGING DIVERGING NOZZLE M=.4388															
1	1														
1C	.01		.4388			1.0	16								
0.0	0.0	0.0	-8.571E-04	-3.995E-02	-3.473E-03	-8.004E-02	-7.978E-03	-1.204E-01							
-1.458E-C2	-1.611E-01	-2.356E-02	-2.021E-01	-3.525E-02	-2.433E-01	-4.997E-02	-2.846E-C1								
0.0	0.0	-8.571E-04	3.995E-02	-3.473E-03	8.004E-02	-7.978E-03	1.204E-01								
-1.458E-C2	1.611E-01	-2.356E-02	2.021E-01	-3.525E-02	2.433E-01	-4.997E-02	2.846E-01								



# STEADY SUPersonic FLOW, SONIC LINE INPUT, UNIFORM MESH

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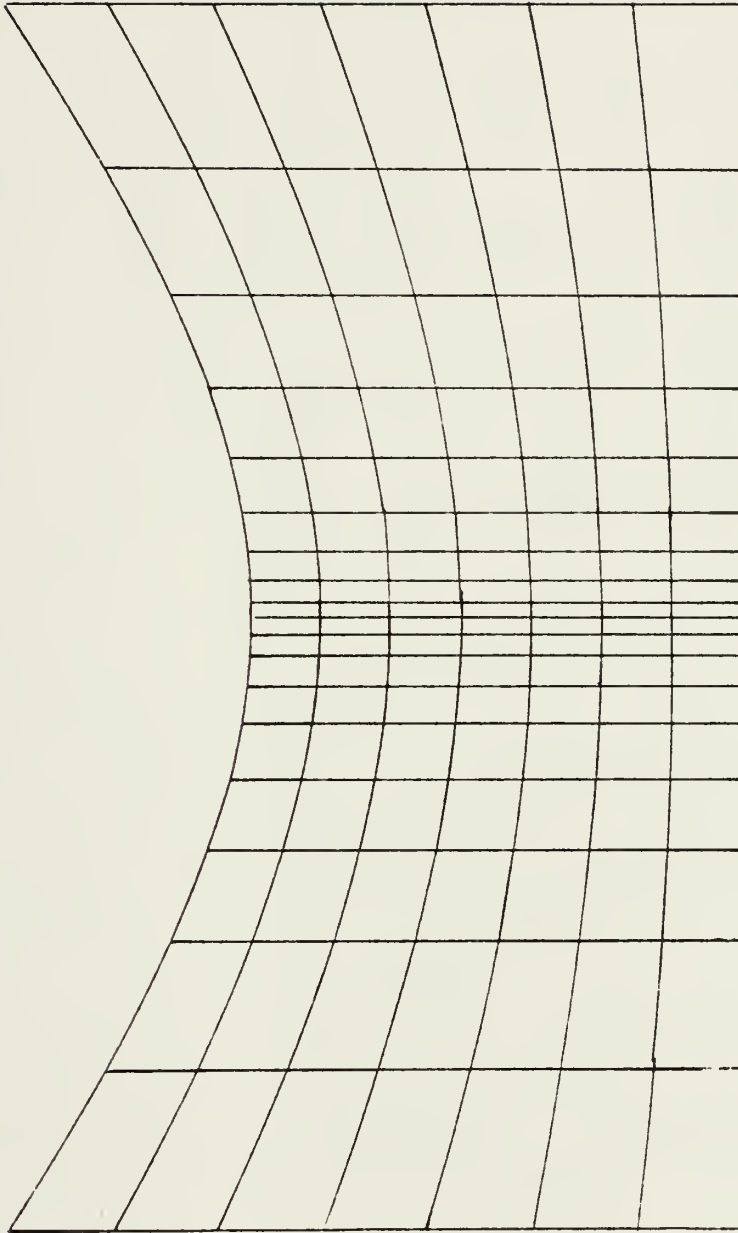


Figure 10 - Finite Element Mesh for the Converging-Diverging Nozzle.





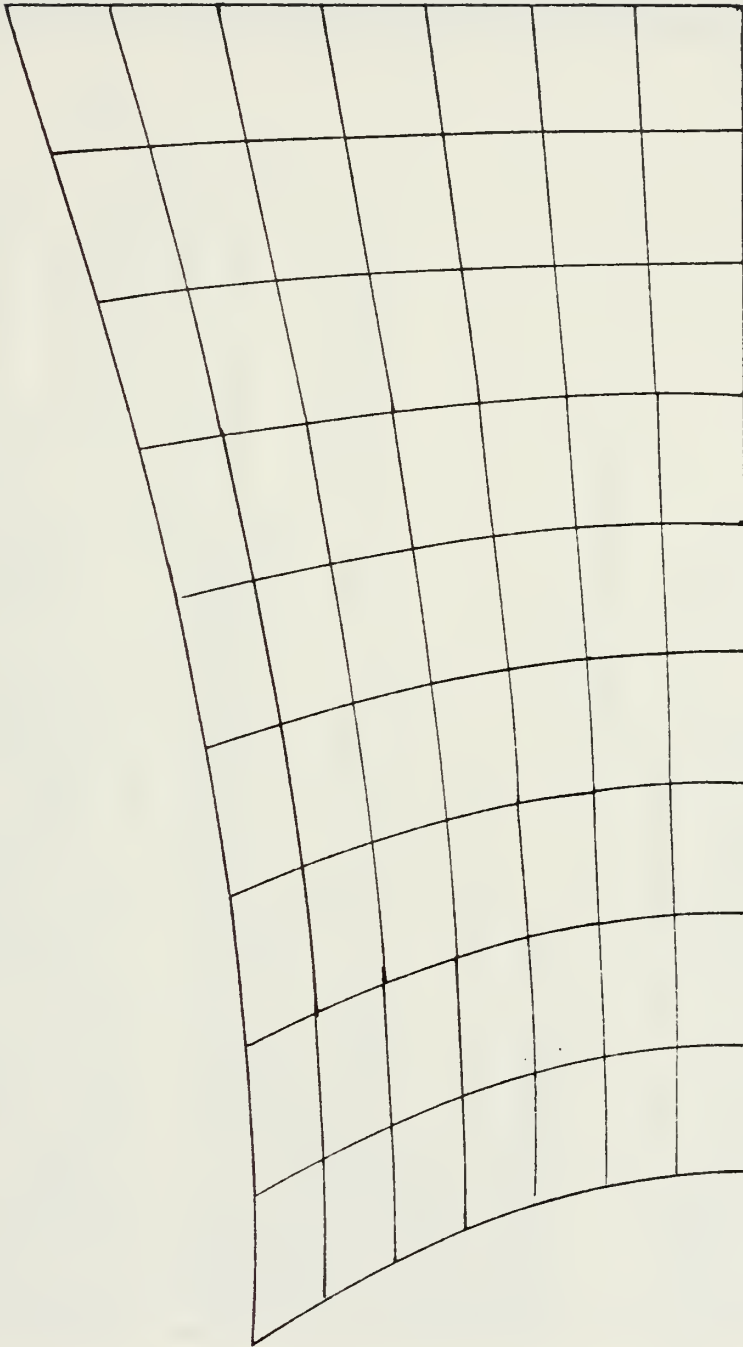


Figure 11 - Finite Element Mesh for the Supersonic Case



## STRANL-II

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STEADY TRANSONIC FLOW ANALYSIS BY FINITE ELEMENT METHOD
USING LEAST SQUARES WITH TRIANGULAR AND QUADRILATERAL ELEMENTS
IOPT(1)=1 USE RESULTS OF PREVIOUS CASE AS STARTING SOLUTION
WHILE THE OTHER OPTION IS IGNORED
IOPT(2)=1, READ IN NON-ZERO INITIAL GUESS
IOPT(3)=1, APPLY LINEARIZED BOUNDARY CONDITIONS ON CHORDLINE
IOPT(4)=1, READ IN NON-ZERO BOUNDARY CONDITIONS
IOPT(5)=1, SONIC LINE VELOCITIES AS INPUTS
THE PROGRAM AS PRESENTLY DIMENSIONED ALLCWS THESE MAXIMA
200 ELEMENTS, 180 NODES, 50 NODES FOR EACH TYPE OF BOUNDARY CCND
MAX FULL BANDWIDTH = 84
DEVELOPED AND CODED BY STEVENS CHAN OF LOCKHEED-HUNTSVILLE ALA.

DIMENSION TITLE(18), IOPT(20), NUD(200,4), S(540,84), SLP(540)
DIMENSION WORK(540,42), UT(540,3), NDATA(3)
DIMENSION X(180), Y(180), RML(180), RMLP(180), CCF(180), NJNT(180)
DIMENSION NIDS(3), NID(3,50), VAF(50), AR(100)
DIMENSION UB(50), VB(50)
LOGICAL LR(50)
EQUIVALENCE(NIDS(1),NFARF),(NIDS(2),NWAKE),(NIDS(3),NBDY)
EQUIVALENCE(NDATA(1),NSK),(NDATA(2),NSM),(NDATA(3),NSC)
DATA NDATA/9,13,14/
DATA NEM,NPM,NCM,NA/200,180,84,50/
NRM=3*NNPM
NFCM=NCM/2
DATA PI/3.1415926/, GAMMA/1.40/
IA=2*NA
CCNST=0.5*(GAMMA-1.0)
EXP=-GAMMA/(GAMMA-1.0)

READ TITLE, CONTROL KEYS AND PROGRAM PARAMETERS

DATA NREAD,NWRITE,NPUNCH/10,8,3/
100 REAC(NREAD,805,END=2000)(TITLE(I),I=1,18)
1C1 CONTINUE
READ(NREAD,820) (IOPT(I),I=1,20)
READ (NREAD,830) ITGIV,ZTEST,RMAC,F1,NFARF$
WRITE (NWRITE,910) (TITLE(I),I=1,18),ZTEST
WRITE (NWRITE,820) (IOPT(I),I=1,20)
WRITE(6,999)ITGIV,ZTEST,RMAC,F1,NFARF$
995 FORMAT(1H0,1ITGIV=,I3,5X,1ZTEST=,F8.4,5X,1RMAC=,F8.4,5X,
11,F1=,F8.4,5X,1NFARF=,I5)
IRES=0
RFT=1.0
SCMAC=RMAC**2

```



```

C
C
C
      IF (IOPT(4) .EQ. 1) READ(NREAD,840) (UB(I), VB(I), I=1, NFARF$)
      IF (IOPT(4) .EQ. 1) CALL BCCND(UB, VB, IOPT(5), NFARF$)
      IF (IOPT(1) .EQ. 1) GO TO 382

      READ AND PRINT MESH DATA, BOUNDARY NODES, AND AIRFCIL SLOPE

      READ (NREAD,825) NELS,NPS,NBW, (NICS(I), I=1,3)
      READ (NREAD,825) ((NOD(I,J), J=1,4), I=1, NELS)
      READ (NREAD,840) (X(I), Y(I), I=1, NPS)
      CC 110 I=1,3
      NS=NIDS(I)
      110 READ (NREAD,825) (NID(I,J), J=1, NS)
      120 READ (NREAD,840) (VAF(I), I=1, NBODY)
      READ (NREAD,825) (NJNT(I), I=1, NPS)

      IF IOPT(3)=1, APPLY LINEARIZED BOUNCARY CCNDITION ON CHORDLINE
      OTHERWISE APPLY NONLINEAR BOUNDARY CONDITIONS ON AIRFOIL SURFACE

      IF (IOPT(3) .NE. 1) GO TO 116
      DO 115 J=1, NBODY
      I=NID(3, J)
      115 Y(I)=0.0
      116 CCNTINUE
      200 WRITE (NWRITE,920) NELS,NPS,NBW
      WRITE (NWRITE,930)
      CC 220 N=1, NELS
      220 WRITE (NWRITE,825) N, (NOD(N, J), J=1,4)
      WRITE (NWRITE,935)
      DO 230 I=1, NPS
      WRITE (NWRITE,940) I, NJNT(I), X(I), Y(I)
      WRITE (NWRITE,951) (NID(1, I), I=1, NFARF)
      WRITE (NWRITE,952) (NID(2, I), I=1, NFAKE)
      WRITE (NWRITE,953) (NID(3, I), I=1, NBODY)
      WRITE (NWRITE,955) (VAF(I), I=1, NBODY)

      REDEFINE MESH DATA, ECT. USING NEW NODAL NUMBERING SYSTEM
C
C
C
      CC 238 N=1, NELS
      DO 238 I=1,4
      IF (NOD(N, I) .EQ.0) GO TO 238
      KK=NOD(N, I)
      NCD(N, I)=NJNT(KK)
      CCNTINUE
      238 DO 239 I=1,3
      IS=NIDS(I)
      DO 239 J=1, IS
      KK=NID(I, J)
      239 NID(I, J)=NJNT(KK)

```

```

TRA000480
TRA000490
TRA000500
TRA000510
TRA000520
TRA000530
TRA000540
TRA000550
TRA000560
TRA000570
TRA000580
TRA000590
TRA000600
TRA000610
TRA000620
TRA000630
TRA000640
TRA000650
TRA000660
TRA000670
TRA000680
TRA000690
TRA000700
TRA000710
TRA000720
TRA000730
TRA000740
TRA000750
TRA000760
TRA000770
TRA000780
TRA000790
TRA000800
TRA000810
TRA000820
TRA000830
TRA000840
TRA000850
TRA000860
TRA000870
TRA000880
TRA000890
TRA000900
TRA000910
TRA000920
TRA000930
TRA000940
TRA000950

```



```

C 243 I=1,NPS
RMLP(I)=X(I)
243 RML(I)=Y(I)
DC 244 II=1,NPS
I=NJNT(II)
X(I)=RMLP(II)
244 Y(I)=RML(II)
C
C HALF BANDWIDTH AND NUMBER OF EQUATIONS
C
NFBW=NBW/2
NEQ=3*NPS
C
C READ NONZERO INITIAL GUESS OR PROCEED WITH ZERO SOLUTION
C
IF (IOPT(2) .NE. 1) GO TO 250
READ (1,840) (S(I,NBW),I=1,NEQ)
GC TO 295
250 DC 260 I=1,NEQ
260 SLP(I)=0.0
C
C ITERATIONS START HERE AND CHECKED IF ITGIV IS EXCEEDED. IF SO,
C PRINT FAIL TO CONVERGE AND PROCEED TO NEXT CASE. OTHERWISE
C CONTINUE TO ITERATE
C
IRES=IRES+1
265 CCNTINUE
IF (IRES .GT. ITGIV) GO TO 600
FORMULATE SYSTEM OF ALGEBRAIC EQUATIONS
DC 266 I=1,NEQ
DC 266 J=1,NBW
266 S(I,J)=0.0
NMAT=1
CALL NEWK(SQMAC,NRM,NCM,NEQ,NBW,NEM,NELS,NOC,SLP,S,CCF,NPM,
1 X,Y,NMAT)
1 IMPCSE B.C. FOR FARFIELD, LINE OF SYMMETRY, AND ON AIRFCIL
IF (IOPT(4) .NE. 0) CALL BNDRY (S,UB,VB,VAF,NID,NFARF,NBODY,
1 NWAKE,NBW,NHBW,NEQ,IOPT,IRES)
1 IF (IOPT(4) .NE. 0) GO TO 276
DC 274 I=1,NFARF
IE=3*NID(1,I)-3
DC 272 II=1,3
IE=IE+1
DC 270 K=1,NBW
S(IE,K)=0.0
270 S(IE,NHBW)=1.0
272 CCNTINUE
274 CCNTINUE
276 CONTINUE

```

TRA00960  
 TRA00970  
 TRA00980  
 TRA00990  
 TRA01000  
 TRA01010  
 TRA01020  
 TRA01030  
 TRA01040  
 TRA01050  
 TRA01060  
 TRA01070  
 TRA01080  
 TRA01090  
 TRA01100  
 TRA01110  
 TRA01120  
 TRA01130  
 TRA01140  
 TRA01150  
 TRA01160  
 TRA01170  
 TRA01180  
 TRA01190  
 TRA01200  
 TRA01210  
 TRA01220  
 TRA01230  
 TRA01240  
 TRA01250  
 TRA01260  
 TRA01270  
 TRA01280  
 TRA01290  
 TRA01300  
 TRA01310  
 TRA01320  
 TRA01330  
 TRA01340  
 TRA01350  
 TRA01360  
 TRA01370  
 TRA01380  
 TRA01390  
 TRA01400  
 TRA01410  
 TRA01420  
 TRA01430





```

DC 280 I=1,NWAKE
IE=3*NID(2,I)
DC 278 K=1,NBW
S(IE,K)=0.0
278 S(IE,NHBW)=1.0
280 DC 285 J=1,NBODY
IE=3*NID(3,J)
DC 282 K=1,NBW
S(IE,K)=0.0
282 S(IE,NHBW)=1.0
IF (IOPT(3) .NE. 1) S(IE,NHBW-1)=-VAF(J)
285 S(IE,NBW)=VAF(J)
CALL BANDED SOLVER TO SOLVE THE SYSTEM OF EQUATIONS
C
C
C
C
STORE STIFFNESS MATRIX AND INITIALIZE STARTING TIME DEPENDENT
UNKNOWN
CALL STCRE(S,WORK,NRM,NCM,NHCM,NFBW,NEQ,NSK)
IF (IRES .GT. 1 .OR. RMAC .GT. .91) GO TO 401
CALL BNCEQ(S,NRM,NCM,NEQ,NHBW)
DC 420 I=1,NRM
UT(I,3)=0.0
420 UT(I,2)=.5*S(I,NBW)
UT(I,1)=S(I,NBW)
GO TO 295
C
ASSEMBLE AND STORE THE DAMPING MATRIX
DC 401 DO 400 I=1,NRM
DC 400 J=1,NCM
400 S(I,J)=0.0
NMAT=2
CALL NEWK(SQMAC,NRM,NCM,NEQ,NBW,NEM,NELS,NOC,SLP,S,CCF,NPM,
1 X,Y,NMAT)
CALL STCRE(S,WORK,NRM,NCM,NHCM,NFBW,NEQ,NSC)
ASSEMBLE AND STORE THE MASS MATRIX
DC 410 I=1,NRM
DC 410 J=1,NCM
410 S(I,J)=0.0
NMAT=3
CALL NEWK(SQMAC,NRM,NCM,NEQ,NBW,NEM,NELS,NOC,SLP,S,CCF,NPM,
1 X,Y,NMAT)
CALL STCRE(S,WORK,NRM,NCM,NHCM,NFBW,NEQ,NSM)
CALL TIME(S,UT,WCRK,NRM,NCM,NHCM,NFBW,NEQ)
C
C
C
PRINT COMPUTED RESULTS
C
C
C
295 WRITE (NWRITE,970) RMAC,RFT
WRITE (NWRITE,975) IRES
CC 305 I=1,NPS

```

TRA01440  
 TRA01450  
 TRA01460  
 TRA01470  
 TRA01480  
 TRA01490  
 TRA01500  
 TRA01510  
 TRA01520  
 TRA01530  
 TRA01540  
 TRA01550  
 TRA01560  
 TRA01580  
 TRA01590  
 TRA01600  
 TRA01610  
 TRA01620  
 TRA01630  
 TRA01640  
 TRA01650  
 TRA01660  
 TRA01670  
 TRA01680  
 TRA01690  
 TRA01700  
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 TRA01970  
 TRA01980  
 TRA01990  
 TRA02000  
 TRA02010  
 TRA02020  
 TRA02030  
 TRA02040  
 TRA02050  
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 TRA02070  
 TRA02080  
 TRA02090  
 TRA02100  
 TRA02110  
 TRA02120  
 TRA02130  
 TRA02140  
 TRA02150  
 TRA02160  
 TRA02170  
 TRA02180  
 TRA02200  
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 TRA02240  
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 TRA02380  
 TRA02390  
 TRA02400  
 TRA02410

```

J=NJNT(I)
II=3*NJNT(I)
PCT=S(II-2,NBW)
UPT=S(II-1,NBW)
V=S(II,NBW)
U=UPT+1.0
QSQ=U*U+V*V
ASQ=CONST*(1.0-QSQ)+1./SQMAC
RML(J)=SQRT(QSQ/ASQ)
PRATIO=(1.0+CONST*RML(J)**2)**EXP
CP=-2.*UPT
COF(J)=SQMAC*(1.0+2.4*UPT)
DELM=RML(J)-RMLP(J)
305 WRITE (NWRITE,978) I,POT,U,V,COF(J),RML(J),FRATIO,CP,DELM
C
C
C   DISPLAY PRESSURE COEFFICIENT CP
C   WRITE (NWRITE,985)
C   ISTOP=0
C   IFNT=0
C   DO 320 J=1,NBODY
C   IF (J.EQ. NBODY) ISTOP =1
C   I=NIID(3,J)
C   LPT=S(3*I-1,NBW)
C   V=S(3*I,NBW)
C   CF=-2.*UPT
C   CCNTINUE
320 IF (IRES .LT. 2) GO TO 382
C
C   CHECK CONVERGENCE. IF SO, PUNCH CCNVERGED SCLUTION AND
C   PROCEED TO NEXT CASE. OTHERWISE, UPDATE SOLUTION AND CONTINUE
C   TC ITERATE
C   DO 340 I=1,NPS
C   PCTE=1.0-RMLP(I)/RML(I)
C   IF (ABS(PCTE) .LT. ZTEST) GO TO 340
C   RFT=FI
C   GO TO 382
340 CCNTINUE
C   WRITE (NPUNCH,840) (S(I,NBW),I=1,NEQ)
C   GO TO 100
382 DO 385 I=1,NPS
385 RMLP(I)=RML(I)
C   RFTC=1.0-RFT
C   CC 390 I=1,NEQ
390 SLP(I)=RFT*S(I,NBW)+RFTC*SLP(I)
C   GO TO 265
600 WRITE (NWRITE,980)

```



```

805 WRITE (1,840) (S(I,NBW),I=1,NEQ)
820 GC TO 100
825 FCRMAT (18A4)
830 FCRMAT (40I2)
835 FCRMAT (16I5)
840 FCRMAT (15,3F10.0,15)
840 FCRMAT (1P8E10.3)
910 FCRMAT (1H1,2X,18A4//, CONVERGENCE LIMIT =',F6.4//)
920 FCRMAT (1H0,NO. OF ELEMENTS=',I4,', NO. CF NODES=',I4,
1 FCRMAT (1H0,FULL BANDWIDTH=',I4//)
930 FCRMAT (1H0,NO. AND ELEMENT NODES//)
935 FCRMAT (1H0,NO. NEW NODE',6X,X(I)',Y(I)')
940 FCRMAT (16,110,2E15.5)
951 FCRMAT (1H0,NO. AT FAR FIELD',(20I5))
952 FCRMAT (1H0,NO. ON THE LINE OF SYMMETRY',(20I5))
953 FCRMAT (1H0,NO. ON THE AIRFOIL',(20I5))
955 FCRMAT (1H0,NO. ON THE AIRFOIL',(8E15.5))
970 FCRMAT (1H1,MACH NUMBER=',F6.3,5X,RELAX. FACTOR =',F6.4//)
975 FCRMAT (1H0,4X,NO. OF ITERATIONS =',I4/
1 FCRMAT (1H0,7X,NO. OF ITERATIONS =',I4/
2 FCRMAT (1H0,7X,NO. OF ITERATIONS =',I4/
980 FCRMAT (1H0,7X,NO. OF ITERATIONS =',I4/
978 FCRMAT (1H0,7X,NO. OF ITERATIONS =',I4/
985 FCRMAT (1H0,7X,NO. OF ITERATIONS =',I4/
2000 STOP
END

```

```

TRA02420
TRA02430
TRA02440
TRA02450
TRA02460
TRA02470
TRA02480
TRA02490
TRA02500
TRA02510
TRA02520
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TRA02590
TRA02600
TRA02610
TRA02620
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TRA02640
TRA02650
TRA02660
TRA02700

```



SUBROUTINE BCOND

```

SUBROUTINE BCOND(U,V,IOPT,N)
DIMENSION U(N),V(N),UP(50),VP(50)
DATA NWRITE/6/
DO 10 I=1,N
  UP(I)=1+U(I)
  VP(I)=V(I)
10 CCNTINUE
  IF (IOPT.EQ.1) GO TO 150
  WRITE(NWRITE,100)
100 FCRMAT( , UCOM AT INLET.)
  N2=N/2
  N3=N2+1
  WRITE(NWRITE,110) (UP(I),I=1,N2)
  WRITE(NWRITE,120) INLET.)
  FCRMAT( , VCOM AT
120 WRITE(NWRITE,110) (VP(I),I=1,N2)
  WRITE(NWRITE,130) EXIT.)
130 FCRMAT( , UCOM AT
  WRITE(NWRITE,140) (UP(I),I=N3,N)
  WRITE(NWRITE,140) EXIT.)
140 FCRMAT( , VCOM AT
110 FCRMAT( , 1H,8(F10.7,5X))
  WRITE(NWRITE,110) (VP(I),I=N3,N)
  GO TO 160
150 WRITE(NWRITE,121)
  WRITE(NWRITE,110) (UP(I),I=1,N)
  WRITE(NWRITE,122)
  WRITE(NWRITE,110) (VP(I),I=1,N)
121 FCRMAT( , UCOM AT SONIC LINE.)
122 FCRMAT( , VCOM AT SONIC LINE.)
160 RETURN
END

```

TRA00010  
 TRA00020  
 TRA00030  
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 TRA00070  
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 TRA00100  
 TRA00110  
 TRA00120  
 TRA00130  
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 TRA00160  
 TRA00170  
 TRA00180  
 TRA00190  
 TRA00200  
 TRA00210  
 TRA00220  
 TRA00230  
 TRA00240  
 TRA00250  
 TRA00260  
 TRA00270  
 TRA00280  
 TRA00290  
 TRA00300  
 TRA00310  
 TRA00320





SUBROUTINE BNDEQ

```

1201 SUBROUTINE BNDEQ(A,NRMAX,NCMAX,N,ITERM)
C      CCNTINUE
C      EQUATION SOLVER FOR BANDED NON-SYMMETRIC SYSTEM OF EQUATIONS
C      SOLUTION STORED IN THE LAST COLUMN AT (I,2*ITERM)
C
C      DIMENSION A(NRMAX,NCMAX)
C      DATA NREAD,NWRITE,NPUNCH/4,8,3/
C      CERO=1.E-6
C      PARE=CERO**2
C      NBND=2*ITERM
C      NEM=NBND-1
C
C      BEGINS ELIMINATION OF THE LOWER LEFT
C
C      DO 1000 I=1,N
C      IF (ABS(A(I,ITERM))) .LT. CERO) GO TO 410
C      GO TO 430
C      IF (ABS(A(I,ITERM))) .LT. PARE) GO TO 1600
C      410 WRITE (6,420) A(I,ITERM), I
C      420 FORMAT (10,420) A(I,ITERM), I
C      430 JLAST=MINO(I+ITERM-1,N)
C      L=ITERM+1
C      CCNTINUE
C      DO 500 J=I,JLAST
C      L=L-1
C      IF (ABS(A(J,L))) .LT. PARE) GO TO 500
C      B=A(J,L)
C      DO 450 K=L,NBND
C      A(J,K)=A(J,K)/B
C      450 IF (I .EQ. N) GO TO 1200
C      CCNTINUE
C      L=0
C      JFIRST=I+1
C      IF (JLAST .LE. I) GO TO 1000
C      DO 900 J= JFIRST,JLAST
C      L=L+1
C      IF (ABS(A(J,ITERM-L))) .LT. PARE) GO TO 900
C      DO 600 K=ITERM,NBM
C      A(J,K-L)=A(J-L,K)-A(J,K-L)
C      600 A(J,NBND)=A(J-L,NBND)-A(J,NBND)
C      IF (I .GE. N-ITERM+1) GO TO 900
C      DO 800 K=1,L
C      A(J,NBND-K)=-A(J,NBND-K)
C      800 CCNTINUE
C      900 CONTINUE
C      1000

```

TRA00010  
 TRA00020  
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 TRA00210  
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 TRA00270  
 TRA00280  
 TRA00290  
 TRA00300  
 TRA00310  
 TRA00320  
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 TRA00360  
 TRA00370  
 TRA00380  
 TRA00390  
 TRA00400  
 TRA00410  
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 TRA00450  
 TRA00460







# SUBROUTINE BNDRY

```

SUBROUTINE BNDRY(S,UB,VB,VAF,NID,NFARF,NBODY,NWAKE,NBW
1,NFBW,NEQ,IOPT,IRES)
DIMENSION S(456,60),UB(50),VB(50),VAF(50),NID(3,50),IOPT(40)
NBW$=NBW-1
IF (IOPT(5) .EQ. 1 .AND. IRES .GT. 1) NFARF=NFARF/2
DC 10 I=1,NFARF
IE=3*NID(1,I)-2
DC 10 J=1,2
IE=IE+1
NX=IE-NHBW
BC=UB(I)
IF (J.EQ.2) BC=VB(I)
DO 41 L=1,NBW$
LX=NX+L
IF (LX .LE. 0 .OR. LX .GT. NEQ) GO TO 41
S(LX,NBW)=S(LX,NBW)-BC*S(LX,NBW-L)
S(LX,(NBW-L))=0.
CCNTINUE
DC 30 K=1,NBW
S(IE,K)=0.0
CCNTINUE
S(IE,NHBW)=1.0
S(IE,NBW)=BC
DC 100 I=1,NBODY
IE=3*NID(3,I)
BC=VAF(I)
NX=IE-NHBW
DC 410 L=1,NBW$
LX=NX+L
IF (LX .LE. 0 .OR. LX .GT. NEQ) GO TO 410
S(LX,NBW)=S(LX,NBW)-BC*S(LX,NBW-L)
S(LX,(NBW-L))=0.0
CCNTINUE
DC 300 K=1,NBW
S(IE,K)=0.0
S(IE,NHBW)=1.0
S(IE,NBW)=BC
DC 200 I=1,NWAKE
IE=3*NID(2,I)
DC 330 K=1,NBW
S(IE,K)=0.0
S(IE,NHBW)=1.
IF (IOPT(5) .EQ. 1 .AND. IRES .GT. 1)NFARF=NFARF*2
RETURN
END

```

TRA00010  
 TRA00020  
 TRA00030  
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 TRA00210  
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 TRA00260  
 TRA00270  
 TRA00280  
 TRA00290  
 TRA00300  
 TRA00310  
 TRA00320  
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 TRA00340  
 TRA00350  
 TRA00360  
 TRA00370  
 TRA00380  
 TRA00390  
 TRA00400  
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 TRA00420  
 TRA00430  
 TRA00440  
 TRA00460



# SUBROUTINE EMTC

```

SUBROUTINE EMTC(A,AT,ATT,XL,YL,PEL,SQMAC)
EVALUATE ELEMENT MATRIX FOR A TRIANGLE BY GAUSSIAN QUADRATURE
SUBROUTINE DERV CALLED TO EVALUATE SHAPE FUNCTION DERIVATIVES
AT THE GAUSSIAN POINTS

DIMENSION A(9,9),P(9),Q(9),NP(5),B(3),C(3),XL(3),YL(3),S(3),
1 DNX(9),DNXX(9),DNY(9),PEL(9),
2 DN(9),AT(9,9),ATT(9,9)
DIMENSION EINT(3,7),WT(7)
DATA LMAX/7/,WT/0.225,3*0.13239415,3*.12593518/
DATA EINT/3*0.33333333,0.05961587,3*0.47C14206,0.05961587,
1 3*0.47C14206,0.05961587,0.79742699,3*0.1C128651,
2 0.79742699,3*0.10128651,0.79742699/
DATA NP/1,2,3,1,2/,GAMMA/1.40/
DATA NREAD,NWRITE,NPUNCH/4,8,3/
DO 2 I=1,9
CC 2 J=1,9
ATT(I,J)=0.
2 A(I,J)=0.0
CC 4 I=1,3
J=NP(I+1)
K=NP(I+2)
B(I)=YL(J)-YL(K)
C(I)=XL(K)-XL(J)
AREA=0.5*(B(2)*C(3)-B(3)*C(2))
CST1=1.0-SQMAC
CST2=SQMAC*(1.0+GAMMA)
DO 10 L=1,LMAX
DO 10 I=1,3
DO 10 J=1,3
S(I)=EINT(I,L)
10 CALL DERV (AREA,B,C,S,DN,DNX,DNXX,DNYY)
U=0.0
UX=0.0
DO 30 I=1,9
U=U+DNX(I)*PEL(I)
30 UX=UX+DNXX(I)*PEL(I)
ALPHA=CST1-CST2*U
DO 40 I=1,9
P(I)=ALPHA*DNXX(I)+DNY(I)
40 Q(I)=P(I)-CST2*UX*DNX(I)
WEIGHT=WT(L)*AREA
CC 60 I=1,9
CST=WEIGHT*Q(I)

```

TRA00010  
 TRA00020  
 TRA00030  
 TRA00040  
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 TRA00060  
 TRA00070  
 TRA00080  
 TRA00090  
 TRA00100  
 TRA00110  
 TRA00120  
 TRA00130  
 TRA00140  
 TRA00150  
 TRA00160  
 TRA00170  
 TRA00180  
 TRA00190  
 TRA00200  
 TRA00210  
 TRA00220  
 TRA00230  
 TRA00240  
 TRA00250  
 TRA00260  
 TRA00270  
 TRA00280  
 TRA00290  
 TRA00300  
 TRA00310  
 TRA00320  
 TRA00330  
 TRA00340  
 TRA00350  
 TRA00360  
 TRA00370  
 TRA00380  
 TRA00390  
 TRA00400  
 TRA00410  
 TRA00420  
 TRA00430  
 TRA00440  
 TRA00450





```

DO 60 J=1,9
  AT(I,J)=AT(I,J)-2*CST*SQMAC*DNX(J)
  ATT(I,J)=ATT(I,J)-CST*SQMAC*DN(J)
  A(I,J)=A(I,J)+CST*P(J)
60 CCNTINUE
1CC RETURN
END

```

```

TRA00460
TRA00470
TRA00480
TRA00490
TRA00500
TRA00510
TRA00530

```



# SUBROUTINE EMQT

```

C
C
C
C
SUBROUTINE EMQT(XQ,YQ,PMQ,SQMAC,EG,EQT,ETT,NTRS)
    GENERATE MATRIX FOR A QUADRILATERAL OR TRIANGLE
    SUBROUTINE EMTC CALLED TO GENERATE MATRIX FOR A BASIC TRIANGLE

    DIMENSION EQ(12,12),ET(9,9),XQ(4),YQ(4),XT(3),YT(3),MP(3,4)
    DIMENSION PMQ(12),PMT(9),EQI(12,12),EQTI(12,12),ETI(9,9),
1    ETIT(9,9)
    DATA MP/1,2,3,3,4,1,2,3,4,4,1,2/
    DATA NREAD,NWRITE,NPUNCH/4,8,3/
    FTOR=1.0
    IF (NTRS.EQ. 4) FTOR=.5
    DO 100 I=1,12
    CC 100 J=1,12
    EQT(I,J)=0.0
    EQTI(I,J)=0.0
    EQ(I,J)=0.0
10C
    DO 150 II=1,NTRS
    CC 105 I=1,3
    NI=MP(I,II)
    IT=3*(I-1)
    IQ=3*(NI-1)
    CO 102 J=1,3
    IT=IT+1
    IC=IQ+1
102
    PMT(IT)=PMQ(IQ)
    XT(I)=XQ(NI)
    YT(I)=YQ(NI)
105
    CALL EMTC(ET,ETT,ETTT,XT,YT,PMT,SQMAC)
    DO 130 K=1,3
    NR=3*(MP(K,II)-1)
    IE=3*(K-1)
    CC 130 KK=1,3
    NR=NR+1
    IE=IE+1
    DO 130 L=1,3
    NC=3*(MP(L,II)-1)
    JE=3*(L-1)
    CC 130 LL=1,3
    NC=NC+1
    JE=JE+1
    EQT(NR,NC)=EQT(NR,NC)+ETT(IE,JE)*FTOR
    EQTI(NR,NC)=EQTI(NR,NC)+ETTT(IE,JE)*FTOR
120
    EC(NR,NC)=EQ(NR,NC)+ET(IE,JE)*FTOR
150
    CC CONTINUE

```

TRA00010  
 TRA00020  
 TRA00030  
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 TRA00070  
 TRA00080  
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 TRA00110  
 TRA00120  
 TRA00130  
 TRA00140  
 TRA00150  
 TRA00160  
 TRA00170  
 TRA00180  
 TRA00190  
 TRA00200  
 TRA00210  
 TRA00220  
 TRA00230  
 TRA00240  
 TRA00250  
 TRA00260  
 TRA00270  
 TRA00280  
 TRA00290  
 TRA00300  
 TRA00310  
 TRA00320  
 TRA00330  
 TRA00340  
 TRA00350  
 TRA00360  
 TRA00370  
 TRA00380  
 TRA00390  
 TRA00400  
 TRA00410  
 TRA00420  
 TRA00430  
 TRA00440  
 TRA00450



RETURN  
END

TRAC0460  
TRA00480



# SUBROUTINE DERV

```

C
C
C
SUBROUTINE DERV(AREA,B,C,S,DN,DNX,DNXX,DNYY)
EVALUATE SHAPE FUNCTION DERIVATIVES AT GAUSSIAN PCINT
DIMENSION B(3),C(3),S(3),DN(9),DNX(9),DNXX(5),DNYY(9),NP(5)
DATA NP/1,2,3,1,2/
CATA NREAD,NWRITE,NPUNCH/4,8,3/
TWCA=2.*AREA
TWCASQ=TWCA**2
CC 200 I=1,3
J=NP(I+1)
K=NP(I+2)
SI=S(I)
SJ=S(J)
SK=S(K)
BI=B(I)
BJ=B(J)
BK=B(K)
CI=C(I)
CJ=C(J)
CK=C(K)
SISQ=SI*SI
BISQ=BI*BI
CISQ=CI*CI
ALFA=0.5*(CK-CJ)
BETA=0.5*(BJ-BK)
F=SI*SJ*SK
FX=BI*SJ*SK+BJ*SK*SI+SJ*BK*BI+SK*BI*BJ)
FXX=2.*(SI*CK*CK+SJ*CK*CI+SK*CI*CJ)
HYY=2.*(SI*(1.0-SI)
CSS=6.*(1.0-2.0*SI)
L=3*I-2
DN(L)=SI*(3-2*SI)+2*H
DNX(L)=BI*CSS+2.*HX
DNXX(L)=BISQ*CS+2.*HXX
DNYY(L)=CISQ*CS+2.*HYY
CS=CK*SJ-CJ*SK
L=L+1
DN(L)=SISQ*CS+ALPHA*H
DNX(L)=2.*BI*SI*CS+TWCA*SI*SQ+ALFA*HX
DNXX(L)=2.*BI*SI*CS+4.*BI*TWCA*SI+ALFA*HXX
DNYY(L)=2.*CISQ*CS+ALFA*HYY
ES=BJ*SK-BK*SJ
L=L+1
TRA00010
TRA00020
TRA00030
TRA00040
TRA00050
TRA00060
TRA00070
TRA00080
TRA00090
TRA00100
TRA00110
TRA00120
TRA00130
TRA00140
TRA00150
TRA00160
TRA00170
TRA00180
TRA00190
TRA00200
TRA00210
TRA00220
TRA00230
TRA00240
TRA00250
TRA00260
TRA00270
TRA00280
TRA00290
TRA00300
TRA00310
TRA00320
TRA00330
TRA00340
TRA00350
TRA00360
TRA00370
TRA00380
TRA00390
TRA00400
TRA00410
TRA00420
TRA00430
TRA00440
TRA00450

```





TRA00460  
TRA00470  
TRA00480  
TRA00490  
TRA00500  
TRA00510  
TRA00520  
TRA00530  
TRA00540  
TRA00550

```

DN(L)=SISQ*BS+BETA*H
CNX(L)=2.*BI*SI*BS+BETA*HX
CNXX(L)=2.*BISQ*BS+BETA*HXX
CNY(L)=2.*CISQ*BS+4.*CI*TWOA*SI+BETA*HYY
CC 300 I=1,9
CNX(I)=DNX(I)/TWOA
DNXX(I)=DNXX(I)/TWOASQ
DNY(I)=DNY(I)/TWOASQ
RETURN
ENC
200
30C

```



# SUBROUTINE NEWK

```

C
C
C
C
C
SUBROUTINE NEWK(SQMAC,NRM,NCM,NEQ,NBW,NEM,NELS,NOD,SLP,S,
1 COEF,NPM,X,Y,NMAT)
GENERATE SYSTEM MATRIX BY ASSEMBLING CONTRIBUTIONS FROM
ALL THE ELEMENTS
SUBROUTINE EMQT CALLED TO GENERATE ELEMENT MATRIX
C
DIMENSION COEF(NPM),X(NPM),Y(NPM),XQ(4),YQ(4),PM(12),BB(12,12)
DIMENSION BBC(12,12),BBM(12,12)
DIMENSION NOD(NEM,4),S(NRM,NCM),SLP(NRM)
DATA NREAD,NWRITE,NPUNCH/4,8,3/
NFEW=NBW/2
CC 480 N=1,NELS
I1=1
IF (NOD(N,4)) 402,402,404
402 NPEL=3
NTRS=1
GC TO 410
404 NPEL=4
NTRS=4
I2=0
CC 408 I=1,4
NI=NOD(N,I)
IF (COEF(NI) .GT. 1.00) I2=I2+1
408 IF (I2.EQ.0) NTRS=2
IF (I2.EQ.4) I1=3
41C DO 425 I=1,NPEL
NI=NOD(N,I)
XQ(I)=X(NI)
YQ(I)=Y(NI)
DO 425 J=1,3
IS=3*(NI-1)+J
IE=3*(I-1)+J
PM(IE)=SLP(IS)
425 CALL EMQT(XQ,YQ,PM,SQMAC,BB,BBC,BEM,NTRS)
DO 450 I=1,NPEL
NFE=3*(NOD(N,I)-1)
IE=3*(I-1)
DO 450 II=1,3
NR=NR+1
IE=IE+1
DC 450 J=1,NPEL
NC=3*(NOD(N,J)-1)-NR+NHBW
JE=3*(J-1)
CC 450 JJ=1,3
TRA000010
TRA000020
TRA000030
TRA000040
TRA000050
TRA000060
TRA000070
TRA000080
TRA000090
TRA000100
TRA000110
TRA000120
TRA000130
TRA000140
TRA000150
TRA000160
TRA000170
TRA000180
TRA000190
TRA000200
TRA000210
TRA000220
TRA000230
TRA000240
TRA000250
TRA000260
TRA000270
TRA000280
TRA000290
TRA000300
TRA000310
TRA000320
TRA000330
TRA000340
TRA000350
TRA000360
TRA000370
TRA000380
TRA000390
TRA000400
TRA000410
TRA000420
TRA000430
TRA000440
TRA000450

```



```

NC=NC+1
JE=JE+1
IF (NMA T-2) 443,442,441
441 S(NR,NC)=S(NR,NC)+BBM(IE,JE)
GO TO 450
442 S(NR,NC)=S(NR,NC)+BBC(IE,JE)
GC TO 450
443 S(NR,NC)=S(NR,NC)+BB(IE,JE)
45C CCNTINUE
48C RETURN
ENC

```

```

TRA00460
TRA00470
TRA00480
TRA00490
TRA00500
TRA00510
TRA00520
TRA00530
TRA00540
TRA00550
TRA00560
TRA00580

```



# SUBROUTINE STORE

```

SUBROUTINE STORE (S,WCRK,NRM,NCM,NF-CM,NHBW,NEQ,NDATA)
DIMENSION S(NRM,NCM),WORK(NRM,NHCM)
REWIND NDATA
1C  DO 10 I=1,NRM
    DO 10 J=1,NHCM
        WCRK(I,J)=0.0
        CC 20 I=1,NEQ
            DO 20 J=1,NHBW
                WCRK(I,J)=S(I,J)
                WRITE (NDATA) WORK
                CC 30 I=1,NEQ
                    DO 30 J=1,NHBW
                        JA=NHBW+J
                        WCRK(I,J)=S(I,JA)
                        WRITE (NDATA) WORK
                    REWIND NDATA
                RETURN
            END
    END

```

STG000010  
 STG000020  
 STG000030  
 STG000040  
 STC000050  
 STC000060  
 STG000070  
 STG000080  
 STG000090  
 STG000100  
 STG000110  
 STC000120  
 STG000130  
 STG000140  
 STC000150  
 STG000160  
 STC000170  
 STG000180





# SUBROUTINE TIME

```

C
C
C
SUBROUTINE TIME(S,U,WORK,NRM,NCM,NFCM,NHBW,NEG)
HCLBOLT INTEGRATION FOR UNSTEADY PROBLEMS
DIMENSION S(NRM,NCM),WORK(NRM,NHCM),U(NRM,3),NDATA(3)
DATA NDATA/9,13,14/
CATA DELT/100./
NBW=2*NHBW
NBW1=NBW-1
AC=2./DELT**2
A1=11./(6*DELT)
A2=5./DELT**2
A3=3./DELT
A4=-2.*A0
A5=-A3/2.
A6=A0/2.
A7=A3/9
CC 100 I=1,NRM
CC 100 J=1,NCM
DO 25 K=1,3
NDATA$=NDATA(K)
REAC (NCATA$) WORK
IF (K-2) 1,2,3
1 CCNST=1
2 CCNST=A0
3 CCNST=A1
4 CCNST=A1
10 S(I,J)=WORK(I,J)*CONST+S(I,J)
C
C
C
MULTIPLY CONTRIBUTING MATRICES BY THE TIME CEPENDENT
VECTORS TO FORM RIGHT HAND SIDE
DO 20 J=1,NHBW
JM=NHBW+1-J
IM=0
CC 20 I=JM,NEG
IM=IM+1
IF (K-2) 11,12,13
11 UW=0.
12 UW=A2*U(IM,1)+A4*U(IM,2)+A6*U(IM,3)
CC TO 20

```

TIM00010  
 TIM00020  
 TIM00030  
 TIM00040  
 TIM00050  
 TIM00060  
 TIM00070  
 TIM00080  
 TIM00090  
 TIM00100  
 TIM00110  
 TIM00120  
 TIM00130  
 TIM00140  
 TIM00150  
 TIM00160  
 TIM00170  
 TIM00180  
 TIM00190  
 TIM00200  
 TIM00210  
 TIM00220  
 TIM00230  
 TIM00240  
 TIM00250  
 TIM00260  
 TIM00270  
 TIM00280  
 TIM00290  
 TIM00300  
 TIM00310  
 TIM00320  
 TIM00330  
 TIM00340  
 TIM00350  
 TIM00360  
 TIM00370  
 TIM00380  
 TIM00390  
 TIM00400  
 TIM00410  
 TIM00420  
 TIM00430  
 TIM00440  
 TIM00450



```

C
12 UM=A3*U(IM,1)+A5*U(IM,2)+A7*U(IM,3)
20 S(IM,NBW)=S(IM,NBW)+WORK(I,J)*UM
25 CCNT=INUE
FCRM UPPER TRIANGULAR EFFECTIVE MATRIX
CC 45 K=1,3
NCATA$=NDATA(K)
READ (NCATA$) WORK
IF(K-2) 21,22,23
21 CCNST=1
CC TO 24
22 CCNST=A0
CC TO 24
23 CCNST=A1
24 CC 30 J=1,NHBW
CC 30 I=1,NEQ
JA=J+NHBW
S(I,JA)=S(I,JA)+WORK(I,J)*CONST
3C MAXJ=NHBW-1
CC 40 J=1,MAXJ
JA=NHBW+J
MAXI=NEQ-J
IM=J
DC 40 I=1,MAXI
IM=IM+1
IF (K-2) 31,32,33
31 UM=0
GC TO 40
32 UM=A2*U(IM,1)+A4*U(IM,2)+A6*U(IM,3)
GC TO 40
33 UM=A3*U(IM,1)+A5*U(IM,2)+A7*U(IM,3)
40 S(IM,NBW)=S(IM,NBW)+WORK(I,J)*UM
45 CCNT=INUE
843 FCRMAT(IH,4(F10.5,5X))
CALL BNDEQ(S,NRM,NCM,NEQ,NHBW)
CC 50 I=1,NRM
U(I,3)=U(I,2)
U(I,2)=U(I,1)
U(I,1)=S(I,NBW)
50 FORMAT(IH,6(F10.5,2X))
8CC RETURN
END

```

```

TIM00460
TIM00470
TIMC0480
TIM00490
TIM00500
TIM00510
TIM00520
TIM00530
TIM00540
TIM00550
TIM00560
TIM00570
TIM00580
TIM00590
TIM00600
TIM00610
TIM00620
TIM00630
TIM00640
TIM00650
TIM00660
TIM00670
TIM00680
TIM00690
TIM00700
TIM00710
TIM00720
TIM00730
TIM00740
TIM00750
TIM00760
TIM00770
TIM00780
TIM00790
TIM00800
TIM00810
TIM00820
TIM00830
TIM00840
TIM00850
TIM00870

```



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tion by the finite ele-  
ment method.

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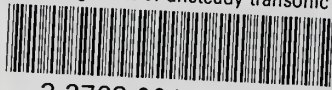
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